

# NEW GENERALIZATION OF FRACTIONAL KINETIC EQUATION USING MULTIVARIABLE ALEPH-FUNCTION

F.Y. AYANT<sup>1</sup>

<sup>1</sup> Teacher in High School , France

**Abstract :** Recently A.Choudhary et al [5] use the multivariable H-function for solving generalized fractional kinetic equation. Motivated by the recent work , we present a new generalization of fractional kinetic equation by using multivariable Aleph-function. The new generalization can be used for the computation of the change of chemical composition in stars like the sun. The solution of the generalized fractional kinetic equation involving multivariable Aleph-function is obtained with help of the Laplace transform method. Further, the same generalized fractional kinetic equation is solved by using the Sumudu transform method. The solution of the proposed problem is presented in a compact form in term of the multivariable Aleph-function. Some special cases, involving the multivariable H-function, the H-function of two variables and the Aleph function of one variable are also considered.

**Keywords :** Multivariable Aleph-function. Fractional Kinetic equation. Laplace transform. Sumudu Transform. Riemann-Liouville fractional integral.

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## 1. INTRODUCTION

The fractional calculus has many important developments and concepts in mathematics initiates with fractional kinetic models (kinetic equation). The great use of Mathematical Physics in imposing astrophysical problems have pulled stargazers and physicists to pay more attention in available mathematical tools that can be used to solve several problems of astrophysics. The importance of fractional kinetic equation has been increased by virtue of its occurrence in certain problem related to kinetic motion of particles in science and engineering. The thermal and hydrostatics equilibrium are pretended as spherically symmetric, non-magnetic, non-rotating, self gravitating model of a star like sun. The properties of star are characterized by its mass, brightness effective surface temperature, radius, central density, temperature etc. Turn over an arbitrary reaction characterized by  $N = N(t)$  which is dependent on time. It is possible to compute rate of change  $dN/dt$  to a balance between the demolition rate  $d$  and the production rate  $p$  of  $N$ , that is  $dN/dt = -d + p$ . In general, through interaction mechanism, demolition and production depend on the quantity  $N$  itself :  $d = d(N)$  or  $p = p(N)$ . This dependence is complicated for the demolition of production at time depends not only on  $N(t)$ , but also on the preceding history  $N(\bar{t})$ ,  $\bar{t} < t$ , of the variable  $N$ . This may be formally represented by [6].

$$\frac{dN}{dt} = -d(N_t) + p(N_t), \quad (1.1)$$

Where  $N_t$  denote the function defined by  $N_t(\bar{t}) = N(t - \bar{t})$ ,  $\bar{t} > 0$ .

Haubold and Mathai [6] studied a special case of this equation, when instance of changes in quantity  $N(t)$  are unvalued, is given by the equation :

$$\frac{dN_i}{dt} = -c_i N_i(t) \quad (1.2)$$

with the initial condition that  $N_i(t = 0) = N_0$  is the number density of species  $i$  at time  $t = 0$ ; constant  $c_i > 0$ , known as standard kinetic equation. The solution of the Eq. (2) is give :

$$N_i(t) = N_0 \exp(-c_i t) \quad (1.3)$$

Alternative form of Eq. (2) can be obtained on integration :

$$N(t) - N_0 = c_0 D_t^{-1} N(t) \quad (1.4)$$

where  ${}_0D_t^{-1}$  is the standard integral operator. Haubold and Mathai [6] have given the fractional generalization of the standard kinetic Eq. (2) as

$$N(t) - N_0 = c_0 D_t^{-\nu} N(t) \quad (1.5)$$

Where  ${}_0D_t^{-\nu}$  is the well known Riemann-Liouville fractional integral operator ( Oldhman and Spanier [10] ; Samko et al [11] ; Miller and Ross [9] , Srivastava and Saxena [14] ) defined by :

$${}_0D_t^{-\nu} = \frac{1}{\Gamma(\nu)} \int_0^t (t-u)^{\nu-1} f(u) du, \quad \text{Re}(\nu) > 0 \quad (1.6)$$

The solution of the fractional Kinetic Eq. ( 6 ) is given by Haubold and Mathai [6] as:

$$N(t) = N_0 \sum_{n=0}^{\infty} \frac{(-1)^n}{\Gamma(\nu n + 1)} (ct)^{n\nu} \quad (1.7)$$

## 2. MATHEMATICAL PREREQUISITES

The multivariable aleph-function is defined in term of multiple Mellin-Barnes type integral :

$$\begin{aligned} \aleph(z_1, \dots, z_r) &= \aleph_{p_i, q_i, \tau_i; R: p_i(1), q_i(1), \tau_i(1); R^{(1)}; \dots; p_i(r), q_i(r), \tau_i(r); R^{(r)}}^{0, n: m_1, n_1, \dots, m_r, n_r} \left( \begin{matrix} z_1 \\ \cdot \\ \cdot \\ \cdot \\ z_r \end{matrix} \right) \\ &= [(a_j; \alpha_j^{(1)}, \dots, \alpha_j^{(r)})_{1, n}]_{\dots} [\tau_i(a_{ji}; \alpha_{ji}^{(1)}, \dots, \alpha_{ji}^{(r)})_{n+1, p_i}] : \\ &\quad \dots [\tau_i(b_{ji}; \beta_{ji}^{(1)}, \dots, \beta_{ji}^{(r)})_{m+1, q_i}] : \\ &\quad \left( \begin{matrix} [(c_j^{(1)}, \gamma_j^{(1)})_{1, n_1}], [\tau_i(c_{ji}^{(1)}, \gamma_{ji}^{(1)})_{n_1+1, p_i^{(1)}}]; \dots; [(c_j^{(r)}, \gamma_j^{(r)})_{1, n_r}], [\tau_i(c_{ji}^{(r)}, \gamma_{ji}^{(r)})_{n_r+1, p_i^{(r)}}] \\ [(d_j^{(1)}, \delta_j^{(1)})_{1, m_1}], [\tau_i(d_{ji}^{(1)}, \delta_{ji}^{(1)})_{m_1+1, q_i^{(1)}}]; \dots; [(d_j^{(r)}, \delta_j^{(r)})_{1, m_r}], [\tau_i(d_{ji}^{(r)}, \delta_{ji}^{(r)})_{m_r+1, q_i^{(r)}}] \end{matrix} \right) \\ &= \frac{1}{(2\pi\omega)^r} \int_{L_1} \dots \int_{L_r} \psi(s_1, \dots, s_r) \prod_{k=1}^r \phi_k(s_k) z_k^{s_k} ds_1 \dots ds_r \quad (2.1) \end{aligned}$$

with  $\omega = \sqrt{-1}$

$$\psi(s_1, \dots, s_r) = \frac{\prod_{j=1}^n \Gamma(1 - a_j + \sum_{k=1}^r \alpha_j^{(k)} s_k)}{\sum_{i=1}^R [\tau_i \prod_{j=n+1}^{p_i} \Gamma(a_{ji} - \sum_{k=1}^r \alpha_{ji}^{(k)} s_k) \prod_{j=1}^{q_i} \Gamma(1 - b_{ji} + \sum_{k=1}^r \beta_{ji}^{(k)} s_k)]} \quad (2.2)$$

$$\text{and } \phi_k(s_k) = \frac{\prod_{j=1}^{m_k} \Gamma(d_j^{(k)} - \delta_j^{(k)} s_k) \prod_{j=1}^{n_k} \Gamma(1 - c_j^{(k)} + \gamma_j^{(k)} s_k)}{\sum_{i(k)=1}^{R^{(k)}} [\tau_{i(k)} \prod_{j=m_k+1}^{p_{i(k)}} \Gamma(1 - d_{ji(k)}^{(k)} + \delta_{ji(k)}^{(k)} s_k) \prod_{j=n_k+1}^{p_{i(k)}} \Gamma(c_{ji(k)}^{(k)} - \gamma_{ji(k)}^{(k)} s_k)]} \quad (2.3)$$

where  $j = 1$  to  $r$  and  $k = 1$  to  $r$

Suppose , as usual , that the parameters

$$a_j, j = 1, \dots, p; b_j, j = 1, \dots, q;$$

$$c_j^{(k)}, j = 1, \dots, n_k; c_{ji^{(k)}}^{(k)}, j = n_k + 1, \dots, p_{i^{(k)}};$$

$$d_j^{(k)}, j = 1, \dots, m_k; d_{ji^{(k)}}^{(k)}, j = m_k + 1, \dots, q_{i^{(k)}};$$

$$\text{with } k = 1 \dots, r, i = 1, \dots, R, i^{(k)} = 1, \dots, R^{(k)}$$

are complex numbers , and the  $\alpha' s, \beta' s, \gamma' s$  and  $\delta' s$  are assumed to be positive real numbers for standardization purpose such that

$$U_i^{(k)} = \sum_{j=1}^n \alpha_j^{(k)} + \tau_i \sum_{j=n+1}^{p_i} \alpha_{ji}^{(k)} + \sum_{j=1}^{n_k} \gamma_j^{(k)} + \tau_{i^{(k)}} \sum_{j=n_k+1}^{p_{i^{(k)}}} \gamma_{ji^{(k)}}^{(k)} - \tau_i \sum_{j=1}^{q_i} \beta_{ji}^{(k)} - \sum_{j=1}^{m_k} \delta_j^{(k)} - \tau_{i^{(k)}} \sum_{j=m_k+1}^{q_{i^{(k)}}} \delta_{ji^{(k)}}^{(k)} \leq 0 \quad (2.4)$$

The reals numbers  $\tau_i$  are positives for  $i = 1$  to  $R$  ,  $\tau_{i^{(k)}}$  are positives for  $i^{(k)} = 1$  to  $R^{(k)}$

The contour  $L_k$  is in the  $s_k$ -p lane and run from  $\sigma - i\infty$  to  $\sigma + i\infty$  where  $\sigma$  is a real number with loop , if necessary , ensure that the poles of  $\Gamma(d_j^{(k)} - \delta_j^{(k)} s_k)$  with  $j = 1$  to  $m_k$  are separated from those of  $\Gamma(1 - a_j + \sum_{i=1}^r \alpha_j^{(i)} s_k)$  with  $j = 1$  to  $n$  and  $\Gamma(1 - c_j^{(k)} + \gamma_j^{(k)} s_k)$  with  $j = 1$  to  $n_k$  to the left of the contour  $L_k$  . The condition for absolute convergence of multiple Mellin-Barnes type contour (1.9) can be obtained by extension of the corresponding conditions for multivariable H-function given by as :

$$|arg z_k| < \frac{1}{2} A_i^{(k)} \pi, \text{ where}$$

$$A_i^{(k)} = \sum_{j=1}^n \alpha_j^{(k)} - \tau_i \sum_{j=n+1}^{p_i} \alpha_{ji}^{(k)} - \tau_i \sum_{j=1}^{q_i} \beta_{ji}^{(k)} + \sum_{j=1}^{n_k} \gamma_j^{(k)} - \tau_{i^{(k)}} \sum_{j=n_k+1}^{p_{i^{(k)}}} \gamma_{ji^{(k)}}^{(k)} + \sum_{j=1}^{m_k} \delta_j^{(k)} - \tau_{i^{(k)}} \sum_{j=m_k+1}^{q_{i^{(k)}}} \delta_{ji^{(k)}}^{(k)} > 0, \text{ with } k = 1 \dots, r, i = 1, \dots, R, i^{(k)} = 1, \dots, R^{(k)} \quad (2.5)$$

The complex numbers  $z_i$  are not zero. Throughout this document , we assume the existence and absolute convergence conditions of the multivariable Aleph-function.

We may establish the the asymptotic expansion in the following convenient form :

$$\aleph(z_1, \dots, z_r) = O(|z_1|^{\alpha_1} \dots |z_r|^{\alpha_r}), \max(|z_1| \dots |z_r|) \rightarrow 0$$

$$\aleph(z_1, \dots, z_r) = O(|z_1|^{\beta_1} \dots |z_r|^{\beta_r}), \min(|z_1| \dots |z_r|) \rightarrow \infty$$

where, with  $k = 1, \dots, r: \alpha_k = \min[Re(d_j^{(k)} / \delta_j^{(k)})], j = 1, \dots, m_k$  and

$$\beta_k = \max[Re((c_j^{(k)} - 1) / \gamma_j^{(k)})], j = 1, \dots, n_k$$

We will use these following notations in this paper

$$U = p_i, q_i, \tau_i; R; V = m_1, n_1; \dots; m_r, n_r \quad (2.6)$$

$$W = p_{i(1)}, q_{i(1)}, \tau_{i(1)}; R^{(1)}, \dots, p_{i(r)}, q_{i(r)}, \tau_{i(r)}; R^{(r)} \quad (2.7)$$

$$A = \{(a_j; \alpha_j^{(1)}, \dots, \alpha_j^{(r)})_{1,n}\}, \{\tau_i(a_{ji}; \alpha_{ji}^{(1)}, \dots, \alpha_{ji}^{(r)})_{n+1,p_i}\} \quad (2.8)$$

$$B = \{\tau_i(b_{ji}; \beta_{ji}^{(1)}, \dots, \beta_{ji}^{(r)})_{m+1,q_i}\} \quad (2.9)$$

$$C = \{(c_j^{(1)}; \gamma_j^{(1)})_{1,n_1}\}, \tau_{i(1)}(c_{ji(1)}^{(1)}; \gamma_{ji(1)}^{(1)})_{n_1+1,p_{i(1)}}\}, \dots, \{(c_j^{(r)}; \gamma_j^{(r)})_{1,n_r}\}, \tau_{i(r)}(c_{ji(r)}^{(r)}; \gamma_{ji(r)}^{(r)})_{n_r+1,p_{i(r)}}\} \quad (2.10)$$

$$D = \{(d_j^{(1)}; \delta_j^{(1)})_{1,m_1}\}, \tau_{i(1)}(d_{ji(1)}^{(1)}; \delta_{ji(1)}^{(1)})_{m_1+1,q_{i(1)}}\}, \dots, \{(d_j^{(r)}; \delta_j^{(r)})_{1,m_r}\}, \tau_{i(r)}(d_{ji(r)}^{(r)}; \delta_{ji(r)}^{(r)})_{m_r+1,q_{i(r)}}\} \quad (2.11)$$

The multivariable Aleph-function write :

$$\aleph(z_1, \dots, z_r) = \aleph_{U:W}^{0,n;V} \left( \begin{array}{c|c} z_1 & A : C \\ \cdot & \cdot \\ \cdot & \cdot \\ \cdot & \cdot \\ z_r & B : D \end{array} \right) \quad (2.12)$$

A new integral transform introduced by Watugala [15], so-called the Sumudu transform. For further details and fundamental properties of the Sumudu transform, see [14, 1, 2, 4]. Let  $A$  be the class of exponentially bounded function

$$f : \mathbb{R} \rightarrow \mathbb{R}, \text{ that is : } |f(t)| < Me^{-t/\tau_1} \text{ if } t \leq 0, \text{ else } |f(t)| > Me^{t/\tau_2}, t \geq 0$$

where  $M, \tau_1$  and  $\tau_2$  are positive real constants,

the Sumudu transform is defined by :

$$G(u) = S[f(t)] = \int_0^\infty f(ut)e^{-t} dt, u \in (-\tau_1, \tau_2) \quad (2.13)$$

The Sumudu transform has been shown to be the theoretical dual of the Laplace transform.

The Riemann-Liouville fractional integral order  $\nu$  is defined by [11, 14],

$${}_0D_t^{-\nu} N(x, t) = \frac{1}{\Gamma(\nu)} \int_0^t (t-u)^{\nu-1} N(x, u) du, \operatorname{Re}(\nu) > 0 \quad (2.14)$$

The Sumudu transform of Riemann-Liouville fractional integral order  $\nu$  is defined as [8]:

$$S\{{}_0D_t^{-\nu} f(t); u\} = u^\nu \bar{f}(u) \quad (2.15)$$

### 3. GENERALIZED FRACTIONAL KINETIC EQUATION

**Lemma 3.1** The Laplace transform of the multivariable Aleph-function as follows

$$L\{t^{\lambda-1}\aleph(z_1 t^\mu, \dots, z_r t^\mu)\} = u^{-\lambda} \aleph_{U_{10}:W}^{0, n+1:V} \left( \begin{matrix} \frac{z_1}{u} \\ \cdot \\ \cdot \\ \cdot \\ \frac{z_r}{u} \end{matrix} \middle| \begin{matrix} (1 - \lambda; \mu, \dots, \mu), A : C \\ \\ \\ B : D \end{matrix} \right) \quad (3.1)$$

where  $\lambda, z_1, \dots, z_r, u \in \mathbb{C}, \operatorname{Re}(u) > 0, \mu > 0, U_{10} = p_i + 1, q_i, \tau_i; R$

**Proof** For convenience, we denote the left side of (3.1) by  $\mathbb{L}$ .

$$\mathbb{L} = \frac{1}{(2\pi\omega)^r} \int_0^\infty \exp(-ut) t^{\lambda-1} \int_{L_1} \dots \int_{L_r} \psi(s_1, \dots, s_r) \prod_{i=1}^r \phi_i(s_i) (z_i t^\mu)^{s_i} ds_1 \dots ds_r dt$$

Changing the order of integration, which permissible under the stated conditions and applied the formula of Laplace transform, we have:

$$\mathbb{L} = \frac{u^{-\lambda}}{(2\pi\omega)^r} \int_{L_1} \dots \int_{L_r} \psi(s_1, \dots, s_r) \prod_{i=1}^r \phi_i(s_i) z_i^{s_i} u^{-\mu(s_1 + \dots + s_r)} \Gamma(\lambda + \mu(s_1 + \dots + s_r)) ds_1 \dots ds_r$$

After simple adjustment we finally arrived at (3.1).

**Lemme 3.2.** From the Lemma3,1, it is clear that

$$L^{-1}\{u^{-\lambda}\aleph(z_1 u^\mu, \dots, z_r u^\mu)\} = t^{\lambda-1} \aleph_{U_{01}:W}^{0, n:V} \left( \begin{matrix} \frac{z_1}{t} \\ \cdot \\ \cdot \\ \cdot \\ \frac{z_r}{t} \end{matrix} \middle| \begin{matrix} A : C \\ \\ (\lambda; \mu, \dots, \mu), B : D \end{matrix} \right) \quad (3.2)$$

where  $\lambda, z_1, \dots, z_r, u \in \mathbb{C}, \operatorname{Re}(u) > 0, \mu > 0, U_{01} = p_i, q_i + 1, \tau_i; R$

**Theorem 3.3.** If  $v > 0, c > 0, d > 0, \mu > 0, \operatorname{Re}(s) > |d|^{v/\alpha}, c \neq d$ , then for the solution of the generalized fractional kinetic equation

$$N(t) - N_0 t^{\mu-1} \aleph(d^v t^v z_0, \dots, d^v t^v z_r) = -c^v {}_0 D_t^{-v} N(t) \quad (3.3)$$

there holds the formula:

$$N(t) = N_0 t^{\mu-1} \sum_{k=0}^{\infty} (-1)^k (ct)^{kv} \aleph_{U_{11}:W}^{0, n+1:V} \left( \begin{matrix} z_1 (dt)^v \\ \cdot \\ \cdot \\ \cdot \\ z_r (dt)^v \end{matrix} \middle| \begin{matrix} (1 - \mu; v \dots v), A : C \\ \\ (1 - kv - \mu; v \dots v), B : D \end{matrix} \right) \quad (3.4)$$

Where  $U_{11} = p_i + 1, q_i + 1, \tau_i; R$

**Proof.** Taking Laplace transform on the both side of (3.3) and using Lemme 3,1, we get

$$\overline{N}(s) - \frac{N_0}{(2\pi\omega)^r} \int_{L_1} \dots \int_{L_r} \psi(s_1, \dots, s_r) \prod_{i=1}^r \phi_i(s_i) \Gamma(\mu + v \sum_{i=1}^r s_i)$$

$$\times s^{-\mu-v} \sum_{i=1}^r s_i \, ds_1 \cdots ds_r = -c^v u^v \overline{N}(s) \quad (3.5)$$

Solving for  $\overline{N}(s)$ , it gives

$$\overline{N}(s) = \frac{N_0}{(2\pi\omega)^r} \int_{L_1} \cdots \int_{L_r} \psi(s_1, \dots, s_r) \prod_{i=1}^r \phi_i(s_i) \Gamma(\mu + v \sum_{i=1}^r s_i) \times s^{-\mu-v} \sum_{i=1}^r s_i \, ds_1 \cdots ds_r \times (1 + c^v u^{-v})^{-1} \quad (3.6)$$

Now, taking inverse Laplace transform on both side of (3.6) and using Lemme 3.2, after little arrangement we finally arrived at the desired result (3.4).

#### 4. ALTERNATIVE METHOD FOR SOLVING GENERALIZED FRACTIONAL KINETIC EQUATION

In this section, we solve the generalized fractional kinetic Eq. (3.3) by an alternate method by using Sumudu transform.

**Lemme 4.1** The Sumudu transform of the multivariable  $\aleph$ -function as follows

$$S\{t^{\lambda-1} \aleph(z_1 t^\mu, \dots, z_r t^\mu)\} = u^{\lambda-1} \aleph_{U_{10}:W}^{0, n+1; V} \left( \begin{matrix} z_1 u^\mu \\ \cdot \\ \cdot \\ \cdot \\ z_r u^\mu \end{matrix} \middle| \begin{matrix} (1 - \lambda; \mu, \dots, \mu), A : C \\ B : D \end{matrix} \right) \quad (4.1)$$

where  $\lambda, z_1, \dots, z_r, u \in \mathbb{C}$ ,  $Re(u) > 0$ ,  $\mu > 0$ ,  $U_{10} = p_i + 1, q_i, \tau_i; R$

**Proof.** In order to prove (4.1) we first express the multivariable  $\aleph$ -function occurring on the left hand side of (4.1) in terms of Mellin-Barnes contour integral and apply Sumudu transform integral (say  $I$ ).

$$\begin{aligned} I &= \int_0^\infty (ut)^{\lambda-1} I(z_1(ut)^\mu, \dots, z_r(ut)^\mu) e^{-t} dt = \\ &= \frac{1}{(2\pi\omega)^r} \int_0^\infty e^{-t} (ut)^{\lambda-1} \int_{L_1} \cdots \int_{L_r} \psi(s_1, \dots, s_r) \prod_{i=1}^r \phi_i(s_i) (z_i(ut)^\mu)^{s_i} ds_1 \cdots ds_r dt \end{aligned}$$

Use the Gamma function  $\int_0^\infty t^{x-1} e^{-t} dt = \Gamma(x)$ , we arrive at

$$I = \frac{u^{\lambda-1}}{(2\pi\omega)^r} \int_{L_1} \cdots \int_{L_r} \psi(s_1, \dots, s_r) \prod_{i=1}^r \phi_i(s_i) z_i^{s_i} \Gamma(\lambda + \mu s_1 + \cdots + \mu s_r) u^{\mu s_1 + \cdots + \mu s_r} ds_1 \cdots ds_r$$

After simple adjustment we finally arrived at (4.1).

**Lemme4.2** From the Lemma 4,1, it is clear that

$$S^{-1}\{u^{\lambda}\aleph(z_1u^{\mu}, \dots, z_ru^{\mu})\} = t^{\lambda-1} = t^{\lambda-1}\aleph_{U_{01}:W}^{0,n:V} \left( \begin{array}{c} z_1t^{\mu} \\ \cdot \\ \cdot \\ \cdot \\ z_rt^{\mu} \end{array} \middle| \begin{array}{c} A : C \\ (\lambda; \mu, \dots, \mu), B : D \end{array} \right) \quad (4.2)$$

where  $\lambda, z_1, \dots, z_r, u \in \mathbb{C}, \operatorname{Re}(u) > 0, \mu > 0, U_{01} = p_i, q_i + 1, \tau_i; R$

**Theorem2.** If  $v > 0, c > 0, d > 0, \mu > 0, \operatorname{Re}(s) > |d|^{v/\alpha}, c \neq d$ , then for the solution of the generalized fractional kinetic equation

$$N(t) - N_0 t^{\mu-1} \aleph(d^v t^v z_0, \dots, d^v t^v z_r) = -c^v {}_0D_t^{-v} N(t) \quad (4.3)$$

there holds the formula :

$$N(t) = N_0 t^{\mu-1} \sum_{k=0}^{\infty} (-1)^k (ct)^{kv} \aleph_{U_{11}:W}^{0,n+1:V} \left( \begin{array}{c} z_1(dt)^v \\ \cdot \\ \cdot \\ \cdot \\ z_r(dt)^v \end{array} \middle| \begin{array}{c} (1 - \mu; v \dots v), A : C \\ (1 - kv - \mu; v \dots v), B : D \end{array} \right) \quad (4.4)$$

Where :  $U_{11} = p_i + 1, q_i + 1, \tau_i; R$

**Proof.** Applying the Sumudu transform on the both side of (4.3) and using Lemme 4.1, we get

$$\begin{aligned} \overline{N}(u) - \frac{N_0}{(2\pi\omega)^r} \int_{L_1} \dots \int_{L_r} \psi(s_1, \dots, s_r) \prod_{i=1}^r (\phi_i(s_i) (d^v z_i)^{s_i}) \Gamma(\mu + v \sum_{i=1}^r s_i) \\ \times u^{\mu+v \sum_{i=1}^r s_i - 1} ds_1 \dots ds_r = -c^v u^{-v} \overline{N}(u) \end{aligned} \quad (4.5)$$

Solving for  $\overline{N}(u)$ , it gives

$$\begin{aligned} \overline{N}(u) = \frac{N_0}{(2\pi\omega)^r} \int_{L_1} \dots \int_{L_r} \psi(s_1, \dots, s_r) \prod_{i=1}^r (\phi_i(s_i) (d^v z_i)^{s_i}) \Gamma(\mu + v \sum_{i=1}^r s_i) \\ \times u^{\mu+v \sum_{i=1}^r s_i - 1} ds_1 \dots ds_r \times (1 + c^v u^{-v})^{-1} \end{aligned} \quad (4.6)$$

Now, taking inverse Sumudu transform both sides of the Eq. (4.6) and using Lemme 4.1, we get after little arrangement the desired result (4.4).

## 5. Particular cases

Remarks : If  $\tau_i = \tau_{i(k)} = 1$ , then the Aleph-function of several variables degenerate in the I-function of several variables defined by Sharma and Ahmad [12].

**Corollary 5.1** If  $v > 0$ ,  $c > 0$ ,  $d > 0$ ,  $\mu > 0$ ,  $Re(s) > |d|^{v/\alpha}$ ,  $c \neq d$ , then for the solution of the generalized fractional kinetic equation

$$N(t) - N_0 t^{\mu-1} I(t^v d^v z_1, \dots, t^v d^v z_r) = -c {}^v_0 D_t^{-v} N(t) \quad (5.1)$$

there holds the formula

$$N(t) = N_0 t^{\mu-1} \sum_{k=0}^{\infty} (-1)^k (ct)^{kv} I_{U_{11}:W}^{0,n+1;V} \left( \begin{matrix} z_1(dt)^v \\ \cdot \\ \cdot \\ z_r(dt)^v \end{matrix} \middle| \begin{matrix} (1 - \mu; v \cdots v), A : C \\ (1 - kv - \mu; v \cdots v), B : D \end{matrix} \right) \quad (5.2)$$

Where :  $U_{11} = p_i + 1, q_i + 1; R$

If  $R = R^{(1)} = \dots, R^{(r)} = 1$ , the multivariable I-function degenerate in the multivariable H-function and we have, see Choudhary et al [5].

**Corollary 5.2** If  $v > 0$ ,  $c > 0$ ,  $d > 0$ ,  $\mu > 0$ ,  $Re(s) > |d|^{v/\alpha}$ ,  $c \neq d$ , then for the solution of the generalized fractional kinetic equation

$$N(t) - N_0 t^{\mu-1} H(t^v d^v z_1, \dots, t^v d^v z_r) = -c {}^v_0 D_t^{-v} N(t) \quad (5.3)$$

there holds the formula

$$N(t) = N_0 t^{\mu-1} \sum_{k=0}^{\infty} (-1)^k (ct)^{kv} H_{p+1,q+1;p_1,q_1;\dots;p_r,q_r}^{0,n+1;m_1,n_1;\dots;m_r,n_r} \left( \begin{matrix} z_1(dt)^v \\ \cdot \\ \cdot \\ z_r(dt)^v \end{matrix} \middle| \begin{matrix} (1 - \mu; v \cdots v), (a_{rj}; \alpha'_{rj}, \dots, \alpha_{rj}^{(r)})_{1,p_r} : (a'_j, \alpha'_j)_{1,p'}; \dots; (a_j^{(r)}, \alpha_j^{(r)})_{1,p^{(r)}} \\ (1 - kv - \mu; v \cdots v), (b_{rj}; \beta'_{rj}, \dots, \beta_{rj}^{(r)})_{1,q_r} : (b'_j, \beta'_j)_{1,q'}; \dots; (b_j^{(r)}, \beta_j^{(r)})_{1,q^{(r)}} \end{matrix} \right) \quad (5.4)$$

If  $r = 2$ , the multivariable H-function reduces into H-function of two variables, for more details, see Srivastava et al [13].

**Corollary 5.3** If  $v > 0$ ,  $c > 0$ ,  $d > 0$ ,  $\mu > 0$ ,  $Re(s) > |d|^{v/\alpha}$ ,  $c \neq d$ , then for the solution of the generalized fractional kinetic equation



$$N(t) - N_0 t^{\mu-1} H(d^v t^v z_1, d^v t^v z_2) = -c^v {}_0D_t^{-v} N(t) \quad (5.5)$$

there holds the formula

$$N(t) = N_0 t^{\mu-1} \sum_{k=0}^{\infty} (-1)^k (ct)^{kv} H_{p+1, q+1; p_1, q_1; p_2, q_2}^{0, n+1; m_1, n_1; m_2, n_2} \left( \begin{matrix} z_1 (dt)^v \\ z_2 (dt)^v \end{matrix} \middle| \begin{matrix} (1-\mu; v, v), (a_j; \alpha'_j, \alpha''_j)_{1,p} \\ (1-kv - \mu; v, v), (b_j; \beta'_j, \beta''_j)_{1,q} \\ (c'_j; \gamma'_j)_{1,p_1}; (c''_j; \gamma''_j)_{1,p_2} \\ (d'_j; \delta'_j)_{1,q_1}; (d''_j; \delta''_j)_{1,q_2} \end{matrix} \right) \quad (5.6)$$

If  $r = 1$ , the multivariable Aleph-function reduces in Aleph-function of one variable and consequently there holds the following result.

**Corollary 5.4** If  $v > 0, c > 0, \mu > 0$ , then the solution of the equation

$$N(t) - N_0 t^{(\mu-1)} \aleph_{p_i, q_i, c_i; r}^{m, n} \left( d^v t^v z \middle| \begin{matrix} (a_j, A_j)_{1, n}, [c_i(a_{ji}, A_{ji})]_{n+1, p_i; r} \\ (b_j, B_j)_{1, m}, [c_i(b_{ji}, B_{ji})]_{m+1, q_i; r} \end{matrix} \right) = -c^v {}_0D_t^{-v} N(t) \quad (5.7)$$

is given :

$$N(t) = N_0 t^{\mu-1} \sum_{k=0}^{\infty} (-1)^k (ct)^{kv} \aleph_{p_i+1, q_i+1, c_i; r}^{m, n+1} \left( d^v t^v z \middle| \begin{matrix} (1-\mu; v, v), \\ (1-kv - \mu; v, v), \\ (a_j, A_j)_{1, n}, [c_i(a_{ji}, A_{ji})]_{n+1, p_i; r} \\ (b_j, B_j)_{1, m}, [c_i(b_{ji}, B_{ji})]_{m+1, q_i; r} \end{matrix} \right) \quad (5.8)$$

For more details ,see ,Kumar et al [7]

## 6. Conclusion and remarks

Multivariable Aleph-function is general in nature and includes a number of known and new results as particular cases. This extended fractional kinetic equation can be used to compute the particle reaction rate and may be utilized in other branch of mathematics. Results obtained in this paper provide an extension of [3, 6, 7, 8].

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Personal adress : 411 Avenue Joseph Raynaud

Le parc Fleuri , Bat B

83140 , Six-Fours les plages

Tel : 06-83-12-49-68

Department : VAR

Counry : FRANCE