# Local Existence and Uniqueness of Solutions to Yang-Mills Heat Flow Problem 

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Abstract - In this paper, we consider the following nonlinear equation

$$
\frac{\partial u}{\partial t}=\partial_{r}^{2} u+\frac{d+1}{r} \partial_{r} u-3(d-2) u^{2}-(d-2) r^{2} u^{3}
$$

where $u:(r, t) \in \mathbb{R}_{+}^{2} \mapsto \mathbb{R}, d \in \mathbb{N}^{*}$. This equation has been investigated by Grotowski in 2001 in studying the Yang-Mills heat flow connections on Riemann manifolds. In the paper, we prove the local Cauchy problem for above equation that is wellposed in $L_{1+r^{2}}^{\infty}\left(\mathbb{R}_{+}\right)$. More precisely, for any initial data $u_{0} \in L_{1+r^{2}}^{\infty}\left(\mathbb{R}_{+}\right)$, there exists $T\left(u_{0}\right)>0$ such that the above equation has a unique solution $u(t) \in L_{1+r^{2}}^{\infty}\left(\mathbb{R}_{+}\right)$for all $t \in\left[0, T\left(u_{0}\right)\right]$.

Keywords - Local Cauchy problem, Local existence and uniqueness problem, Yang-Mills heat flows, Yang-Mills connections, Geometric flows.

## I. INTRODUCTION

In this paper, we are interested in the following nonlinear heat equation

$$
\left\{\begin{array}{l}
\frac{\partial u}{\partial t}=\partial_{r}^{2} u+\frac{d+1}{r} \partial_{r} u-3(d-2) u^{2}-(d-2) r^{2} u^{3},(r, t) \in(0,+\infty) \times[0, T),  \tag{1}\\
u(0)=u_{0} \in L^{\infty}\left(\mathbb{R}_{+}\right),
\end{array}\right.
$$

where $u:(r, t) \in \mathbb{R}_{+}^{2} \mapsto \mathbb{R}, d \in \mathbb{N}^{*}$. Equation (1) has a strong connection to the problem of Yang Mills connections in $\mathbb{R}^{d} \times S O(d)$ which is an important part in Yang-Mills theory. We would like to mention that the theory was used to study the weak nuclear forces, governing the nuclear decay of some particles which is considered as a non-commutative version of Maxwell's electromagnetism, see more details in [1]-[6]. In particular, the Yang-Mills heat flow has received a lot of attention from both mathematics and physics communities. Results on existence and uniqueness of weak solutions in other functional spaces were obtained in [7] for $d=2,3$, and [8] and [9] for $d=4$; global existence has proved in [10]-[12]; the regularity was established in [13] in for higher dimensions; stability were proved in [14] and [15]; singularity formations have been studied in [16]-[24], the local well posedness was studied in [25] and [26] in Sobolev spaces . In particular, the study of YangMills connections usually is considered in abstract spaces such as Riemann manifolds which get lot of inconveniences. For that reason, the author in [17] reduced the study of Yang-Mills connections to the problem introduced in (1).

The main goal of this paper is to study the local Cauchy problem to (1) since strong connections to the Yang-Mills problem. Recently, the authors in [24] solved the local Cauchy problem for (1) in $C_{0}^{\infty}\left(\mathbb{R}_{+}\right)$defined as the set of smooth functions with compact supports. However, it remains open if the local Cauchy problem in can be solved $L^{\infty}\left(\mathbb{R}_{+}\right)$. In this work, we aim to prove the Cauchy problem in $L_{1+r^{2}}^{\infty}\left(\mathbb{R}_{+}\right) \subset L^{\infty}\left(\mathbb{R}_{+}\right)$which is strictly larger than $C_{0}^{\infty}\left(\mathbb{R}_{+}\right)$.

For more convenience, we mention below some important notations which will be used in the proof of the paper. Let $\Omega$ be a Lebesgue measurable set in $\mathbb{R}^{n}, n \in \mathbb{N}^{*}$, we denote $L^{\infty}(\Omega)$ as the set of all Lebesgue measurable functions $f$ on $\Omega$ satisfying

$$
\inf \{B \text { such that } \mu(\{x \in \Omega:|f(x)|>B\})=0\}<+\infty
$$

which is a Banach space with the following norm

$$
\|f\|_{L^{\infty}(\Omega)}=\inf \{B \text { such that } \mu(\{x \in \Omega:|f(x)|>B\})=0\} .
$$

Similarly, we also define $L_{1+|x|^{\infty}}^{\infty}(\Omega)$

$$
L_{1+|x|^{2}}^{\infty}(\Omega)=\left\{f \in L^{\infty}(\Omega) \text { such that }\left\|\left(1+|x|^{2}\right) f(.)\right\|_{L^{\infty}(\Omega)}<+\infty\right\}
$$

and it is also a Banach space with the norm

$$
\|f\|_{L_{1+1+)^{\infty}}^{\infty}(\Omega)}=\left\|\left(1+|x|^{2}\right) f(.)\right\|_{L^{\infty}(\Omega)}
$$

We call $f$ a radially symmetric function on $\Omega$ if and only if for all orthogonal matrices $A$ and $x \in \Omega$, then it satisfies that $A x \in \Omega$ and $f(A x)=f(x)$. Hence, we introduce

$$
L_{\text {rad }}^{\infty}(\Omega)=\left\{f \in L^{\infty}(\Omega) \text { such that } f \text { radially symmetric }\right\},
$$

which is a Banach space with $\|\cdot\|_{L^{\infty}(\Omega)}$ norm. In particular, we also define the abstract space of Banach-valued functions $X^{T}, T>0$ by

$$
X^{T}=L^{\infty}\left([0, T], L_{r a d}^{\infty}(\Omega) \cap L_{1+|x|^{2}}^{\infty}(\Omega)\right)
$$

which is also a Banach space with the norm

$$
\|z\|_{X^{T}}=\|Z(t)\|_{L^{\infty}([0, T])} \text { where } Z(t)=\|z(t)\|_{L_{1+t \times]^{\infty}}^{\infty}(\Omega)}
$$

Let $\Delta$ be Laplace operator in Euclide space $\mathbb{R}^{d+2}$ defined by $\Delta=\sum_{j=1}^{d+2} \partial_{x_{j} x_{j}}^{2}$. By taking $\Omega=\mathbb{R}^{d+2}$, we recall the definition of semi-group $\left\{e^{\Delta t}\right\}_{t>0}$ as follows

$$
e^{\Delta t}: L^{\infty}\left(\mathbb{R}^{d+2}\right) \rightarrow L^{\infty}\left(\mathbb{R}^{d+2}\right)
$$

and

$$
\begin{equation*}
e^{\Delta t} f(x)=\frac{1}{(4 \pi t)^{\frac{d+2}{2}}} \int_{\mathbb{R}^{d+2}} f(y) e^{-\frac{|x-y|^{2}}{4 t}} d y \tag{2}
\end{equation*}
$$

We would like to mention [27, Proposition 48.4] the following fundamental estimate

$$
\begin{equation*}
\left\|e^{\Delta t} f\right\|_{L^{\infty}\left(\mathbb{R}^{d+2}\right)} \leq\|f\|_{L^{\infty}\left(\mathrm{x}^{d+2}\right)} \tag{3}
\end{equation*}
$$

## II. MAIN RESULTS

In this section, we aim to prove the local Cauchy problem for equation (1) in $L_{1+|x|^{2}}^{\infty}\left(\mathbb{R}_{+}\right)$. However, it is so hard to give a direct proof to equation (1) due to complexity of the linear operator

$$
\partial_{r}^{2}+\frac{d+1}{r} \partial_{r} .
$$

To overcome this challenge, we used the idea investigated in [28] where the authors successfully handled the Cauchy problem for harmonic heat flows. Let $f$ be a function defined in $\mathbb{R}_{+}$, then we denote $\tilde{f}$ as $f$ 's extension on $\mathbb{R}^{d+2}$ given by

$$
\begin{equation*}
\tilde{f}(x)=f(|x|) \tag{4}
\end{equation*}
$$

We can see that the extension is always a radially symmetric function on $\mathbb{R}^{d+2}$. In particular, once $f \in C^{2}\left(\mathbb{R}_{+}\right)$, then we have $\tilde{f} \in C^{2}\left(\mathbb{R}^{d+2}\right)$ and the following identity holds

$$
\Delta \tilde{f}(x)=\partial_{r}^{2} f+\frac{d+1}{r} \partial_{r} f(|x|)
$$

Let us consider $u$ is a $C^{2}\left(\mathbb{R}_{+}\right)$-solution to (1) (so-called the classical solution), then the extension $\tilde{u}$ belongs to $C^{2}\left(\mathbb{R}^{d+2}\right)$ and $\tilde{u}$ satisfies the following equation

$$
\left\{\begin{array}{l}
\partial_{t} \tilde{u}=\Delta \tilde{u}-3(d-2) \tilde{u}^{2}-(d-2)|x|^{2} \tilde{u}^{3},(x, t) \in \mathbb{R}^{d+2} \times(0, T)  \tag{5}\\
\tilde{u}(0)=\tilde{u}_{0} \in L_{r a d}^{\infty}\left(\mathbb{R}^{d+2}\right)
\end{array}\right.
$$

From the symmetricity of equation (5), the solution $\tilde{u}$ remains radially symmetric as well as the solution exists. In the following, we aim to prove the local Cauchy problem to equation (5) in $L_{r a d}^{\infty}\left(\mathbb{R}^{d+2}\right) \cap L_{1+|x|^{\infty}}^{\infty}\left(\mathbb{R}^{d+2}\right)$. Firstly, we have the following Lemma:
Lemma 2.1. Let us consider $\alpha$ be a positive number and $f \in L^{\infty}\left(\mathbb{R}^{d+2}\right)$ satisfying

$$
\begin{equation*}
\left\|\left(1+|x|^{\alpha}\right) f(.)\right\|_{L^{\infty}\left(\mathbb{R}^{d+2}\right)} \leq A \tag{6}
\end{equation*}
$$

Then, it holds that

$$
\left\|\left(1+|x|^{\alpha}\right) e^{t \Delta} f\right\|_{L^{\infty}\left(\mathbb{R}^{d+2}\right)} \leq C(\alpha) A, t>0
$$

where the semi-group $e^{t \Delta}$ defined as in (2).
Proof: By the explicit formula in (2), we write as follows

$$
e^{\Delta t} f(x)=\frac{1}{(4 \pi t)^{\frac{d+2}{2}}} \int_{\mathbb{R}^{d+2}} e^{-\frac{|x-y|^{2}}{4 t}} f(y) d y
$$

Let us consider $R>0$ large enough and fixed later. For all $|x| \leq R$, we use (6) to derive

$$
\left|\left(1+|x|^{\alpha}\right) e^{\Delta t} f(x)\right| \leq \frac{A\left(1+R^{\alpha}\right)}{(4 \pi t)^{\frac{d+2}{2}}} \int_{\mathbb{R}^{d+2}} \frac{e^{-\frac{|x-y|^{2}}{4 t}}}{1+|y|^{\alpha}} d y \leq C(R, \alpha) A
$$

It remains to the case $|x| \geq R$. Let us define $K_{1}=\left\{y:|y| \leq \frac{|x|}{2}\right\}$ and $K_{1}=\left\{y:|y|>\frac{|x|}{2}\right\}$. It is obvious that $K_{1} \cup K_{2}=\mathbb{R}^{d+2}$.
Then, we have the decomposition

$$
\left(1+|x|^{\alpha}\right) e^{\Delta t} f(x)=\frac{1}{(4 \pi t)^{\frac{d+2}{2}}} \int_{\mathbb{R}^{\alpha+2}} \frac{1+|x|^{\alpha}}{1+|y|^{\alpha}} e^{-\frac{|x-y|^{2}}{4 t}}\left(1+|y|^{\alpha}\right) f(y) d y=I_{1}+I_{2}
$$

where

$$
I_{1}=\frac{1}{(4 \pi t)^{\frac{d+2}{2}}} \int_{K_{1}} \frac{1+|x|^{\alpha}}{1+|y|^{\alpha}} e^{-\frac{|x-y|^{2}}{4 t}}\left(1+|y|^{\alpha}\right) f(y) d y
$$

and

$$
I_{2}=\frac{1}{(4 \pi t)^{\frac{d+2}{2}}} \int_{K_{2}} \frac{1+|x|^{\alpha}}{1+|y|^{\alpha}} e^{-\frac{|x-y|^{2}}{4 t}}\left(1+|y|^{\alpha}\right) f(y) d y
$$

+ For $I_{1}:$ We use the facts that $t \in(0,1)$ and $y \in K_{1},|x-y| \geq|x|-|y| \geq \frac{|x|}{2}$, then it follows

$$
\left|I_{1}\right| \leq \frac{A}{(4 \pi t)^{\frac{d+2}{2}}} \int_{K_{1}}\left(1+|x|^{\alpha}\right) e^{-\frac{|x-y|^{2}}{4 t}} d y
$$

Using the changing variable $z=\frac{x-y}{2 \sqrt{t}}$, we get the following

$$
\begin{aligned}
& \frac{1}{(4 \pi t)^{\frac{d+2}{2}}} \int_{K_{1}}\left(1+|x|^{\alpha}\right) e^{-\frac{|x-y|^{2}}{4 t}} d y \\
& =\frac{\left(1+|x|^{\alpha}\right)}{(\pi)^{\frac{d+2}{2}}} \int_{y \in K_{1}, z=\sqrt{y} y-x} e^{-z^{2}} d z \\
& \leq \frac{\left(1+|x|^{\alpha}\right)}{(\pi)^{\frac{d+2}{2}}} \int_{\frac{\mid x}{2} \leq|x| \frac{3|x|}{2}} e^{-z^{2}} d z \leq C(1+|x|)^{d+2+\alpha} e^{-\frac{|x|^{2}}{4}} \\
& \leq C(R, \alpha),
\end{aligned}
$$

which implies $\left|I_{1}\right| \leq C(R, \alpha) A$.

+ For $I_{2}$ : Using the fact that $y \in K_{2},|y| \geq \frac{|x|}{2}$, then it follows $\frac{1+|x|^{\alpha}}{1+|y|^{\alpha}} \leq C(\alpha)$. Hence, we estimate $I_{2}$ as follows

$$
\left|I_{2}\right| \leq \frac{C(\alpha) A}{(4 \pi t)^{\frac{d+2}{2}}} \int_{K_{2}} e^{-\frac{|x-y|^{2}}{4 t}} d y \leq C(\alpha) A
$$

Thus, we conclude that for all $|x| \geq R$, we have

$$
\left|\left(1+|x|^{\alpha}\right) e^{\Delta t} f(x)\right| \leq C(R, \alpha) A .
$$

Finally, we get the conclusion of the proof.
Now, we have the following result:
Proposition 2.2. Let $\tilde{u}_{0} \in L_{\text {rad }}^{\infty}\left(\mathbb{R}^{d+2}\right) \cap L_{1+\mid x x^{2}}^{\infty}\left(\mathbb{R}^{d+2}\right)$ arbitrarily, there exists $T=T\left(\tilde{u}_{0}\right)>0$ such that problem (5) has a unique solution on $[0, T]$ and $\tilde{u}(t) \in L_{\text {rad }}^{\infty}\left(\mathbb{R}^{d+2}\right) \cap L_{1+|x|^{\infty}}^{\infty}\left(\mathbb{R}^{d+2}\right)$ for all $t \in[0, T]$. In particular, the following estimate holds

$$
\|\tilde{u}\|_{X^{r}} \leq\left\|\tilde{u}_{0}\right\|_{L^{L+1+p^{\infty}}\left(\mathbb{x}^{d+2}\right)}+1 .
$$

Proof: The result follows from the Banach fixed point theorem (the unique fixed point of a constructive mapping). Let us consider initial data $\tilde{u}_{0} \in L_{\text {rad }}^{\infty}\left(\mathbb{R}^{d+2}\right) \cap L_{1+\mid x x^{2}}^{\infty}\left(\mathbb{R}^{d+2}\right)$ and $T>0$ and define

$$
B_{0}=\left\{\tilde{u} \in X^{T} \text { such that }\left\|\tilde{u}-e^{t \Delta}\left(\tilde{u}_{0}\right)\right\|_{X^{T}} \leq 1\right\} .
$$

In addition, we define $\Gamma$ a mapping on $X^{T}$

$$
\begin{equation*}
\Gamma(\tilde{u})(t)=e^{t \Delta} \tilde{u}_{0}+\int_{0}^{t} e^{(t-s) \Delta} F(s, \tilde{u}(s)) d s, t>0 \text { and } \Gamma(z)(0)=\tilde{u}_{0} \tag{7}
\end{equation*}
$$

and $F$ defined by

$$
\begin{equation*}
F(\tilde{u})=-3(d-2) \tilde{u}^{2}-(d-2)|x|^{2} \tilde{u}^{3} . \tag{8}
\end{equation*}
$$

In the below, we aim to prove that once $T$ small enough, $\Gamma$ satisfies the following properties:
(H1): $\Gamma$ maps $B_{0}$ into itself.
(H2): $\Gamma$ is a contraction mapping on $B_{0}$ i.e. there exists $\lambda \in(0,1)$ such that

$$
\left\|\Gamma\left(\tilde{u}_{1}\right)-\Gamma\left(\tilde{u}_{2}\right)\right\|_{x^{r}} \leq \lambda\left\|\tilde{u}_{1}-\tilde{u}_{2}\right\|_{x^{T}}, \forall \tilde{u}_{1}, \tilde{u}_{2} \in B_{0} .
$$

- Proof of (H1): Let $\tilde{u} \in B_{0}$ arbitrarily, we derive from (8) and the fact $\|\tilde{u}\|_{x^{T}} \leq 1$

Applying Lemma 2.1 with $\alpha=2$ and $t-s \in(0,1)$, we get

$$
\left\|e^{(t-s) \Delta} F(\tilde{u}(s))\right\|_{L_{1+4 x^{2}}^{\infty}\left(\mathbb{x}^{d+2}\right)} \leq C_{1}\|\tilde{u}\|_{X^{T}}
$$

which yields

$$
\begin{equation*}
\left\|\Gamma(\tilde{u})-e^{t \Delta} \tilde{u}_{0}\right\|_{X^{T}}=\left\|\int_{0}^{t} \int^{(t-s) \Delta} F(\tilde{u}(s)) d s\right\|_{X^{T}} \leq C_{1} \sqrt{T}\|\tilde{u}\|_{X^{T}}, \tag{9}
\end{equation*}
$$

for all $\tilde{u} \in B_{0}$. Taking $T \leq\left(\frac{1}{C_{1}}\right)^{2}$, we get

$$
\left\|\Gamma(\tilde{u})-e^{t s} \tilde{u_{0}}\right\|_{X^{T}} \leq\|\tilde{u}\|_{X^{T}}, \forall \tilde{u} \in B_{0} .
$$

In addition, since $\tilde{u}$ is radially symmetric, so $F(\tilde{u})$ defined as in (8) is and the convolution in (2) saves the symmetry that leads $\Gamma(\tilde{u}(t))$ is radially symmetric for all $t \in[0, T]$. Finally, we conclude (H1).

- Proof of (H2): Let $\tilde{u}_{1}, \tilde{u}_{2} \in B_{0}$, and we write

$$
\begin{aligned}
F\left(\tilde{u}_{1}(s)\right)-F\left(\tilde{u}_{2}(s)\right) & =-3(d-2)\left(\tilde{u}_{1}(s)-\tilde{u}_{2}(s)\right)\left(\tilde{u}_{1}(s)+\tilde{u}_{2}(s)\right) \\
& -(d-2)|x|^{2}\left(\tilde{u}_{1}(s)-\tilde{u}_{2}(s)\right)\left(\tilde{u}_{1}^{2}(s)+\tilde{u}_{1}(s) \tilde{u}_{2}(s)+\tilde{u}_{2}^{2}(s)\right) .
\end{aligned}
$$

Since $\left\|\tilde{u}_{1}\right\|_{X^{r}} \leq 1$ and $\left\|\tilde{u}_{2}\right\|_{X^{x}} \leq 1$, we derive

Regarding to Lemma 2.1, we have

$$
\left\|e^{(t-s) \Delta}\left(F\left(\tilde{u}_{1}(s)\right)-F\left(\tilde{u}_{2}(s)\right)\right)\right\|_{L_{1++p^{\infty}}^{\infty}\left(\mathbb{x}^{d+2}\right)} \leq C\left\|\tilde{u}_{1}(s)-\tilde{u}_{2}(s)\right\|_{L_{1+1+\alpha^{2}}^{\infty}\left(\mathbb{x}^{d+2}\right)} \leq C\left\|\tilde{u}_{1}-\tilde{u}_{2}\right\|_{X^{T}},
$$

for all $t-s>0, s \in(0, t)$. Thus, we derive from (6) that for all $t \in(0, T]$

$$
\left\|\Gamma\left(\tilde{u}_{1}\right)(t)-\Gamma\left(\tilde{u}_{2}\right)(t)\right\|_{L_{1+1+x^{2}}{ }^{2}\left(\mathbb{R}^{d+2}\right)} \leq C\left\|\tilde{u}_{1}-\tilde{u}_{2}\right\|_{X^{T}}\left(\int_{0}^{t} d s\right) \leq C_{2} T\left\|\tilde{u}_{1}-\tilde{u}_{2}\right\|_{X^{T}},
$$

which yields

$$
\left\|\Gamma\left(\tilde{u}_{1}\right)-\Gamma\left(\tilde{u}_{2}\right)\right\|_{x^{T}} \leq C_{2} T\left\|\tilde{u}_{1}-\tilde{u}_{2}\right\|_{X^{\tau}}, \forall \tilde{u}_{1}, \tilde{u}_{2} \in B_{0} .
$$

Finally, we choose $C_{2} \sqrt{T} \leq \frac{1}{2}$ i.e. $T \leq\left(\frac{1}{2 C_{2}}\right)^{2}$ then

$$
\left\|\Gamma\left(\tilde{u}_{1}\right)-\Gamma\left(\tilde{u}_{2}\right)\right\|_{X^{r}} \leq \frac{1}{2}\left\|\tilde{u}_{1}-\tilde{u}_{2}\right\|_{X^{r}}, \forall \tilde{u}_{1}, \tilde{u}_{2} \in B_{0},
$$

which concludes (H2).
Now, we continue on the proof of the proposition, since $X^{T}$ is a Banach space and $\Gamma$ is a contractive map from $B_{0}$ to itself, so we apply Banach fixed point theorem that there uniquely exists $\tilde{u} \in B_{0}$ such that

$$
\tilde{u}(t)=\Gamma(\tilde{u})(t), \forall t \in[0, T],
$$

and we have the estimate

$$
\begin{equation*}
\|\tilde{u}\|_{X^{T}} \leq\left\|e^{t \Delta} \tilde{u}_{0}\right\|_{X^{T}}+\left\|\Gamma(\tilde{u})-e^{t \Delta} \tilde{u}_{0}\right\|_{X^{T}} \leq\left\|\tilde{u}_{0}\right\|_{L_{1+\mid x x^{2}}^{\infty}\left(\mathbb{R}^{d+2}\right)}+1 \tag{10}
\end{equation*}
$$

In particular, by the parabolic regularity of the semi-group $e^{t \Delta}$, we improve that $\tilde{u}(t) \in C^{2}\left(\mathbb{R}^{d+2}\right), \forall t \in(0, T)$ and it satisfies equation (5) for all $(x, t) \in \mathbb{R}^{d+2} \times(0, T)$ point-wise, and we derive from (10)

$$
|\tilde{u}(x, t)| \leq \frac{\left\|\tilde{u}_{0}\right\|_{L_{1++x]^{\infty}}^{\infty}}+1}{1+|x|^{2}}, \forall(x, t) \in \mathbb{R}^{d+2} \times(0, T]
$$

Finally, we get the conclusion of the Proposition.
Consequently, Proposition 2.2 implies the following result:
Proposition 2.3: Let $d \geq 1$ be an integer number and initial choice $u_{0} \in L_{1+r^{2}}^{\infty}\left(\mathbb{R}_{+}\right)$. Then, there exists $T=T\left(u_{0}\right)$ such that problem (1) has a unique solution on $[0, T]$ and $u(t) \in L_{1+|x|^{2}}^{\infty}\left(\mathbb{R}_{+}\right) \cap C^{2}\left(\mathbb{R}_{+}\right), \forall t \in(0, T]$. In particular, we have the following estimate

$$
|u(r, t)| \leq \frac{C\left(u_{0}\right)}{1+r^{2}}, \forall(r, t) \in \mathbb{R}_{+} \times(0, T)
$$

Proof: Let $u_{0} \in L_{1+r^{2}}^{\infty}\left(\mathbb{R}_{+}\right)$, then, it follows from the extension (4) that $\tilde{u}_{0} \in L_{r a d}^{\infty}\left(\mathbb{R}^{d+2}\right) \cap L_{1+|x|^{2}}^{\infty}\left(\mathbb{R}^{d+2}\right)$. Applying Proposition 2.2, we obtain the existence and the uniqueness of the solution $\tilde{u}$ to equation (5) on $\left[0, T\left(u_{0}\right)\right]$ and $\tilde{u}(t) \in L_{\text {rad }}^{\infty}\left(\mathbb{R}_{+}\right) \cap L_{1+|x|^{2}}^{\infty}\left(\mathbb{R}_{+}\right) \cap C^{2}\left(\mathbb{R}_{+}\right), \forall t \in(0, T]$, then, the problem (1) and (8) are equivalent in the radially symmetric setting, this leads to the existence and the uniqueness of $u$ and also the conclusion the proposition completely follows.

Remark 2.4: We can repeat the proof of Propositions 2.2 and 2.3, to establish the local existence and uniqueness in spaces $L_{1+r^{\alpha}}^{\infty}\left(\mathbb{R}_{+}\right)$where $\alpha \geq \frac{2}{3}$. Since the main difficulty is to handle the hugeness of nonlinear term $r^{2} u^{3}$ at infinity. However, once $u \in L_{1+r^{\alpha}}^{\infty}\left(\mathbb{R}_{+}\right), \alpha \geq \frac{2}{3}$ it follows that the nonlinearity is controlled well by

$$
r^{2} u^{3}=\frac{r^{2}}{\left(1+r^{\alpha}\right)^{3}}\left(\left(1+r^{\alpha}\right) u\right)^{3} \leq\|u\|_{L_{1+r^{\alpha}}^{\infty}}^{3}\left(1+r^{\alpha}\right)^{\frac{2}{\alpha}-3}
$$

since $\frac{2}{\alpha}-3 \leq 0$. We kindly refer the readers to check the details of the general results.

## III. CONCLUSION

As we showed in Proposition 2.3, the local Cauchy problem in $L_{1+r^{2}}^{\infty}\left(\mathbb{R}_{+}\right)$is completely solved (also in $L_{1+r^{\alpha}}^{\infty}\left(\mathbb{R}_{+}\right), \alpha \geq \frac{2}{3}$ as in Remark 2.4). In comparison to the result proved in Donninger and Schörkhuber, 2019 where the authors considered the problem in $C_{0}^{\infty}\left(\mathbb{R}_{+}\right)$, our result is better. The technique of the proof replies to the idea given by Biernat and Seki, 2019 that extended the original problem to the radially symmetric one in $\mathbb{R}^{d+2}$, and then, we established a new property that the semigroup $e^{t \Delta}, t>0$ reserves the polynomial decays showed in Lemma 2.1, then the existence and uniqueness follows by the route map based on Banach fixed point theorem.

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