

Original Article

Positive Periodic Solutions for a Sixth-Order Variable Coefficient Singular Differential Equation with Indefinite Weights

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Abstract - In this paper, we consider a sixth-order variable coefficient singular differential equation with indefinite weights

$$u^{(6)}(t) + \sum_{i=0}^5 a_i(t)u^i(t) = \frac{h(t)}{u^\rho} + e(t),$$

where ω is a positive constant, $h(t), e(t) \in L^1(\mathbb{R}/\omega\mathbb{Z})$ and $e(t) > 0$ for any $t \in [0, \omega]$. By using the Krasnoselskii-Guo fixed point theorem, we prove the existence of positive periodic solutions to the above equation.

Keywords - Positive periodic solutions, Singular differential equations, Green function, Sixth-order differential equation, Indefinite weights.

1. Introduction

Singular differential equations have a wide range of applications, such as physics, engineering, biology, and economics (Chu 2018; Huang et al. 2014) [1],[9]. In recent years, it has been widely applied in various disciplines such as boundary layer theory, Bose-Einstein condensates, and optical signal transmission, and some application examples can be found in (Torres 2015) [19][19][19].

The study of singular differential equations began in 1987, when Lazer and Solimini [15] discussed the existence of positive periodic solutions of the following two types of singular differential equations

$$u'' - \frac{1}{u^\lambda} = \varphi(t) \tag{1.1}$$

and

$$u'' + \frac{1}{u^\lambda} = \varphi(t) \tag{1.2}$$

where λ is a positive real constant and $\varphi(t)$ is a periodic function with period ω . They gave a necessary and sufficient condition for the existence of periodic solutions of equations (1.1) and (1.2).

Generally, for a damped differential equation with singularities

$$u'' + a(t)u = f(t, u) + h(t)$$

where $h(t)$ is an external force and the nonlinearity $f(t, u)$ may be singular at the origin, it is said that it has an attractive at the origin, if

$$\lim_{x \rightarrow 0^+} f(t, u) = -\infty, \forall t \in \mathbb{R}.$$



and has a repulsive singularity at the origin, if

$$\lim_{x \rightarrow 0^+} f(t, u) = +\infty, \forall t \in \mathbb{R}.$$

In 2010, Hakl and Torres [10] investigated the following second-order differential equation with attractive-repulsive singularities

$$u'' = \frac{g(t)}{u^{k_1}} - \frac{h(t)}{u^{k_2}} + \varphi(t). \tag{1.3}$$

The authors obtained the existence of periodic solutions of equation (1.3) by the method of lower and upper solutions.

In 2017, Lu et al. [13] discussed the existence of positive periodic solutions for the following Liénard equation with a singularity of repulsive type

$$u'' + f(u)u' - \frac{a(t)}{u} = \varphi(t)$$

where $a(t)$ is a continuous function with periodic T and $a(t) > 0$.

In 2020, Liu et al. [12] investigated the existence of a positive ω -periodic solution for the equation with a singularity of attractive type

$$u'' + a(t)u = \frac{h(t)}{u^\rho} + e(t),$$

where $a(t), h(t) \in C(\mathbb{R} \setminus \omega\mathbb{Z})$ were negative.

At the same time, more and more scholars have turned their attention to the research of higher-order singular differential equations.

In 2006, Chu and Zhou [6] studied the existence of positive periodic solutions for the third-order singular differential equation

$$u''' + k^3u = f(t, u(t)), 0 \leq t \leq 2\pi, \tag{1.4}$$

where k is a constant, and $f(t, u(t))$ has a singularity at $u = 0$. They proved that equation (1.4) has a positive periodic solution when $k \in (0, \frac{1}{\sqrt{3}})$. In 2010, Li [11] studied the existence of positive periodic solutions of the following third-order singular differential equation

$$u''' = f(t, u(t), u'(t), u''(t)), \tag{1.5}$$

where $f \in C(\mathbb{R} \times (0, +\infty) \times \mathbb{R} \times \mathbb{R}, \mathbb{R})$ is a periodic function of t and has a singularity at $u = 0$, they proved that equation (1.5) has at least one positive periodic solution by using the point theorem in cones.

In 2003, Conti, Terracini, and Verzini [5] studied the fourth-order equation

$$u^{(4)}(t) - cu''(t) = f(t, u(t)), \quad t \in [0, T],$$

with periodic boundary conditions. Given any positive integer $n \geq 1$, they proved there exists a T -periodic solution to the problem with precisely $2n$ simple zeroes in $[0, T]$ by Nehari's argument of combining variational methods and nodal properties of solutions. In 2009, Cui and Zou [7] studied the following singular fourth-order boundary value problems:

$$\begin{cases} u^{(4)}(t) = f(t, u(t), -u''(t)), & 0 < t < 1 \\ u(0) = u(1) = u''(0) = u''(1) = 0 \end{cases}$$

where $f(t, u, y)$ may be singular at $t = 0, u = 0, y = 0$ and $t = 1, u = 0, y = 0$. They proved the existence and uniqueness of

solutions, which are established by constructing a special cone and using the cone compression and expansion fixed-point theorem. In 2015, Cheng and Ren [3] established the existence of positive periodic solutions for the following fourth-order singular nonlinear differential equation with variable-coefficient

$$u^{(4)} + m(t)u''' + n(t)u'' + p(t)u' + q(t)u = h(t, u) + \varphi(t) \tag{1.6}$$

At the same time, Cheng and Ren [4] calculated Green’s functions of the above fourth-order singular nonlinear differential equation in ten cases in detail. In 2017, Xin, Han, and Cheng [15] investigated a class of fourth-order singular nonlinear differential equations with superlinearity or sublinearity assumptions

$$u^{(4)}(t) + au'''(t) + bu''(t) + cu' + du(t) = \lambda g(t)f(u(t)) + \lambda e(t)$$

with $\lambda > 0$ is a positive parameter, and $e(t)$ may take a positive value or a negative value. By employing Green’s function in [4] and Krasnoselskii-Guo fixed point theorem, they proved the existence of positive periodic solutions for the singular fourth-order differential equation.

In 2023, Liu et al. [14] investigated the sixth-order singular differential equation

$$u^{(6)}(t) + \sum_{i=0}^5 a_i(t)u^i(t) = f(t, u) + e(t), \tag{1.7}$$

where $a_i(t)(i = 0, 1, \dots, 5)$ are ω -periodic functions, $e(t) \in L^p([0, \omega])$, $f(t, u) \in C(\mathbb{R} \times \mathbb{R}^+, \mathbb{R})$ is an L^2 -Caratheodory function and $f(t, u)$ is an ω -periodic function in t . The authors calculated the Green functions in 36 cases in detail. They proved the existence of positive solutions to equation (1.7) with attractive singularity or repulsive singularity by using Schauder's fixed-point theorem.

Motivated by [14], in this paper, we employ the result of Green’s function of a sixth-order singular differential equation and consider the following sixth-order singular differential equation with indefinite weights

$$u^{(6)}(t) + \sum_{i=0}^5 a_i(t)u^i(t) = \frac{h(t)}{u^\rho} + e(t) \tag{1.8}$$

where $a_i(t) \in C(\mathbb{R}, \mathbb{R}^+)(i=0, 1, \dots, 5)$ are ω -periodic functions, $h(t), e(t) \in L^1(\mathbb{R}/\omega\mathbb{Z})$ and $e(t) > 0$ for any $t \in [0, \omega]$.

Theorem 1.1. Suppose equations (3.1), (4.2), (4.3), and (4.5) in section 3 hold, then equation (1.8) has at least one positive periodic solution.

The rest of the paper is organized as follows. In Section 2, Green’s function and some useful properties for Green’s function are given. In Section 3, we prove the existence of positive periodic solutions for equation (1.8) by the Krasnoselskii-Guo fixed point theorem. Another linear Green’s function and useful properties are given in the appendix.

2. Green’s Function of a Sixth-Order Linear Differential Equation

In this paper, to study the existence of positive periodic solutions for the sixth-order singular differential equation (1.8), we employ the Green’s functions and their properties of the equation

$$\begin{cases} u^{(6)}(t) + \sum_{i=0}^5 a_i(t)u^i(t) = \varphi(t), \\ u^{(i)}(0) = u^{(i)}(\omega), i = 0, 1, 2, 3, 4, 5, \end{cases} \tag{2.1}$$

where $a_i(t) \in C(\mathbb{R}, \mathbb{R}^+)(i=0, 1, \dots, 5)$, $\varphi(t) \in C(\mathbb{R}, \mathbb{R}^+)$ are ω -periodic functions, which can be found in [14]. The authors reduce and decompose the Green’s functions into 36 cases, and only 6 typical cases are analyzed in detail due to the similarities among these cases. We will take Case 1 as a representative example to discuss the positive periodic solutions of equation (1.8), and Cases 2-6 can be obtained similarly, which will be presented in the appendix.

In this section, we only provide detailed properties of Green’s function of Case 1 as follows :

Case 1: There exist two ω -periodic differentiable functions $\alpha(t), \beta(t): \mathbb{R} \rightarrow \mathbb{R}^+$ with $\int_0^\omega \alpha(s)ds > 0$ and $\int_0^\omega \beta(s)ds > 0$, and

an ω -periodic function $m(t)$ with $m(t) > 0$, such that

$$\begin{aligned} a_5(t) &= 2\alpha(t), \\ a_4(t) &= m(t) + \alpha^2(t) + 2\alpha'(t) + 2\beta(t), \\ a_3(t) &= 4m'(t) + 2m(t)\alpha(t) + 2\beta(t)\alpha(t) + \beta'(t)\alpha(t) + \alpha''(t) + 2\beta'(t), \\ a_2(t) &= 6m''(t) + 6m'(t)\alpha(t) + m(t)[2\alpha'(t) + 2\beta(t) + \alpha^2(t)] + \beta''(t) + \alpha'(t)\alpha(t) + \beta^2(t), \\ a_1(t) &= 4m'''(t) + 6m''(t)\alpha(t) + 2m'(t)[2\alpha'(t) + 2\beta(t) + \alpha^2(t)] + m(t)[\alpha''(t) + 2\beta'(t) + 2\beta(t)\alpha(t)], \\ a_0(t) &= m^{(4)}(t) + m'''(t)[2\alpha(t) + \alpha'(t)\alpha(t)] + m''(t)[2\alpha'(t) + 2\beta(t) + \alpha^2(t)] + m'(t)[2\beta'(t) + 2\alpha(t)\beta(t) + \alpha''(t) + \alpha'(t)\alpha(t)] + m(t)[\beta''(t) + \alpha(t)\beta'(t) + \beta^2(t)]. \end{aligned}$$

In this case, equation (2.1) can be transformed into

$$\begin{cases} v''(t) + \alpha(t)v'(t) + \beta(t)v(t) = \varphi(t), \\ v(0) = v(\omega), v'(0) = v'(\omega), \end{cases}$$

$$\begin{cases} y''(t) + \alpha(t)y'(t) + \beta(t)y(t) = v(t) \\ y(0) = y(\omega), y'(0) = y'(\omega) \end{cases}$$

and

$$\begin{cases} u''(t) + m(t)u(t) = y(t), \\ u(0) = u(\omega), u'(0) = u'(\omega), \end{cases}$$

Since the solutions of these second-order equations can be respectively expressed as

$$v(t) = \int_0^\omega G_1(t, s)\varphi(s)ds, y(t) = \int_0^\omega G_1(t, s)v(s)ds, u(t) = \int_0^\omega G_2(t, s)y(s)ds,$$

Then the equation (2.1) can be written as

$$\begin{aligned} u(t) &= \int_0^\omega G_1(t, \gamma_1) \int_0^\omega G_1(\gamma_1, \gamma_2) \int_0^\omega G_1(\gamma_2, s)\varphi(s)dsd\gamma_2 d\gamma_1 \\ &= \int_0^\omega \int_0^\omega \int_0^\omega G_2(t, \gamma_1) G_1(\gamma_1, \gamma_2)G_1(\gamma_2, s)\varphi(s) dsd\gamma_2d\gamma_1 \\ &= \int_0^\omega \left[\int_0^\omega \int_0^\omega G_2(t, \gamma_1) G_1(\gamma_1, \gamma_2)G_1(\gamma_2, s)d\gamma_2d\gamma_1 \right] \varphi(s) ds, \end{aligned}$$

or

$$u(t) = \int_0^\omega G^1(t, s)\varphi(s)ds,$$

where

$$G^1(t, s) = \int_0^\omega \int_0^\omega G_2(t, \gamma_1)G_1(\gamma_1, \gamma_2)G_1(\gamma_2, s) d\gamma_2d\gamma_1.$$

Lemma 2.1.([14]) Suppose that equations (4.2), (4.3), and (4.5) hold, then $G^1(t, s) > 0$ for any $(t, s) \in [0, \omega] \times [0, \omega]$.

3. The Existence of Positive Periodic Solutions of the Equation

In this section, we first give Krasnoselskii-Guo's fixed-point theorem.

Lemma 3.1.([8]) (Krasnoselskii-Guo fixed point theorem) Let X be a Banach space and $K \subset X$ is a cone. Suppose that Ω_1 and Ω_2 are open subsets of X with $0 \in \Omega_1, \overline{\Omega_1} \subset \Omega_2$. Let

$$\Psi: K \cap (\overline{\Omega_2} \setminus \Omega_1) \rightarrow K$$

be a completely continuous operator such that one of the following conditions holds:

- (i) $\|\Psi u\| \leq \|u\|$ for $u \in K \cap \partial\Omega_1$ and $\|\Psi u\| \geq \|u\|$ for $u \in K \cap \partial\Omega_2$;
- (ii) $\|\Psi u\| \geq \|u\|$ for $u \in K \cap \partial\Omega_1$ and $\|\Psi u\| \leq \|u\|$ for $u \in K \cap \partial\Omega_2$;

Then Ψ has a fixed point in $K \cap (\overline{\Omega_2} \setminus \Omega_1)$.

Now, we consider the existence of positive periodic solutions for equation (1.8) by the Krasnoselskii-Guo fixed point theorem. Due to similarity, we only investigate in detail the existence of Case 1 described in Section 2.

Obviously, $u(t)$ is the unique ω -periodic solution of the equation (2.1) and

$$u(t) = \int_0^\omega G^1(t, s)\varphi(s)ds.$$

Define

$$K := \left\{ u \in C_\omega : \min_{t \in \mathbb{R}} u(t) \geq \sigma \|u\|_\infty \right\},$$

where

$$C_\omega := \{ u \in C(\mathbb{R}) : u(t + \omega) = u(t), \forall t \in \mathbb{R} \}$$

and

$$\|u\|_\infty = \max_{t \in \mathbb{R}} |u(t)|.$$

Obviously, K is a cone of C_ω .

Besides, for any arbitrarily given ω -periodic function $h(t)$ provided we define

$$h^+(t) := \max\{0, h(t)\}, \quad h^-(t) := -\min\{0, h(t)\}, \quad \bar{h} := \frac{1}{\omega} \int_0^\omega h(t)dt.$$

Theorem 3.1. Suppose that $h(t), e(t) \in L^1(\mathbb{R}/\omega\mathbb{Z})$ and $e(t) > 0$ for any $t \in [0, \omega]$. Assume equations (4.2), (4.3), (4.5) and

$$\bar{h}_+ > \frac{1}{A\omega\sigma^{1+\rho}} \left(\frac{\|h^-\|_\infty}{e_*} \right)^{\frac{1+\rho}{\rho}} \tag{3.1}$$

hold, where

$$A := \min_{0 \leq s, t \leq \omega} G^1(t, s), B := \max_{0 \leq s, t \leq \omega} G^1(t, s), \quad \sigma = \frac{A}{B},$$

Then equation (1.8) has one positive ω -periodic solution at least.

Proof: Define

$$(\Psi u)(t) := \int_0^\omega G^1(t, s) \left(\frac{h(s)}{u^\rho} + e(s) \right) ds.$$

It is easy to prove that one ω -periodic solution of equation (1.8) is a fixed point of the operator equation $\Psi u = u$. And from Lemma 2.1, we can obtain $G^1(t, s) > 0$ for any $(t, s) \in [0, \omega] \times [0, \omega]$.

Define two open sets

$$\Omega_1 := \{ u \in C_\omega : \|u\| < r \}, \Omega_2 := \{ u \in C_\omega : \|u\| < R \},$$

where r and R are two positive constants, and from equation (3.1) we obtain

$$R > r = (A\omega\bar{h}^+)^{\frac{1}{1+\rho}} > \frac{1}{\sigma} \left(\frac{\|h^-\|_\infty}{e_*} \right)^{\frac{1}{\rho}}.$$

First, we prove $\Psi(K \cap (\overline{\Omega_2} \setminus \Omega_1)) \subset K$. Actually, for any $u \in K \cap (\overline{\Omega_2} \setminus \Omega_1)$, we have

$$\sigma r \leq u(t) \leq R, \quad \forall t \in \mathbb{R}.$$

Since $r > \frac{1}{\sigma} \left(\frac{\|h^-\|_\infty}{e_*} \right)^{\frac{1}{\rho}}$ we can obtain

$$\begin{aligned} \frac{h(t)}{u^\rho} + e(t) &= \frac{h^+(t)}{u^\rho} - \frac{h^-(t)}{u^\rho} + e(t) \\ &> -\frac{h^-(t)}{u^\rho} + e_* \\ &> -\frac{\|h^-\|_\infty}{(\sigma r)^\rho} + e_* \\ &> 0. \end{aligned} \tag{3.2}$$

From equation (3.2) and $\sigma = \frac{A}{B}$, we get

$$\begin{aligned} \min_{t \in \mathbb{R}} (\Psi u)(t) &= \min_{t \in \mathbb{R}} \int_0^\omega G^1(t, s) \left(\frac{h(s)}{u^\rho} + e(s) \right) ds \\ &\geq A \int_0^\omega \left(\frac{h(s)}{u^\rho} + e(s) \right) ds \\ &= \sigma B \int_0^\omega \left(\frac{h(s)}{u^\rho} + e(s) \right) ds \\ &\geq \sigma \max_{t \in \mathbb{R}} \int_0^\omega G^1(t, s) \left(\frac{h(s)}{u^\rho} + e(s) \right) ds \\ &= \sigma \|\Psi u\|_\infty. \end{aligned}$$

Therefore, $\Psi(K \cap (\overline{\Omega_2} \setminus \Omega_1)) \subset K$. And from the Arzela-Ascoli Theorem, it is easy to prove $\Psi(K \cap (\overline{\Omega_2} \setminus \Omega_1)) \subset K \rightarrow K$ is a compact operator.

Next, we prove

$$\|\Psi u\| \leq \|u\|, \quad \forall u \in K \cap \partial\Omega_2. \tag{3.3}$$

Actually, for any $u \in K \cap \partial\Omega_2$, we can get

$$\|u\| = R \text{ and } \sigma R \leq u(t) \leq R, \forall t \in \mathbb{R}.$$

From $\sigma = \frac{A}{B}$ and $e(t) > 0$ for any $t \in [0, \omega]$, we obtain

$$\begin{aligned} (\Psi u)(t) &= \int_0^\omega G^1(t, s) \left(\frac{h(s)}{u^\rho} + e(s) \right) ds \\ &= \int_0^\omega G^1(t, s) \left(\frac{h^+(s)}{u^\rho} - \frac{h^-(s)}{u^\rho} + e(s) \right) ds \\ &\leq B\omega \left(\frac{\bar{h}_+}{(\sigma R)^\rho} + \bar{e} \right). \end{aligned}$$

Obviously, we can select a suitably large R that satisfies.

$$B\omega \left(\frac{\bar{h}_+}{(\sigma R)^\rho} + \bar{e} \right) \leq R.$$

Thus, we obtain $\|\Psi u\| \leq \|u\|$, which means equation (3.3) holds.

Lastly, we will prove

$$\|\Psi u\| \geq \|u\|, \forall u \in K \cap \partial\Omega_1, \quad (3.4)$$

In fact, for any $u \in K \cap \partial\Omega_1$ We can obviously obtain

$$\|u\| = r \text{ and } \sigma r \leq u(t) \leq r, \forall t \in \mathbb{R}.$$

From $\sigma = \frac{A}{B}$, we obtain

$$\begin{aligned} (\Psi u)(t) &= \int_0^\omega G^1(t, s) \left(\frac{h^+(s)}{u^\rho} - \frac{h^-(s)}{u^\rho} + e(s) \right) ds \\ &> \int_0^\omega G^1(t, s) \left(\frac{h^+(s)}{u^\rho} \right) ds \\ &\geq \frac{A\omega \bar{h}_+}{r^\rho} = r. \end{aligned}$$

where $r = (A\omega \bar{h}_+)^{\frac{1}{1+\rho}}$. Therefore, equation (3.4) holds.

From Lemma 3.1, we can conclude that Q has a fixed point $u \in K \cap (\bar{\Omega}_2 \setminus \Omega_1)$. Obviously, this point is a positive ω -periodic solution of the equation. (1.8) and $u \in [\sigma r, R]$.

Remark 3.1. The existence of positive periodic solutions for equation (1.8) with indefinite weights for Cases 2 and 6 can be shown in a similar way. More precisely, it can be shown for Case 2 by establishing (4.2), (4.3), and substituting (4.5) with $\rho < \frac{2\pi}{\sqrt{3}\omega}$, for Case 3 by establishing (4.5) and replacing (4.2), (4.3) with $\rho < \frac{2\pi}{\sqrt{3}\omega}$, for Case 4 by establishing (4.2), (4.3), for Case 5 by establishing (4.5), and for Case 6 by replacing (4.2), (4.3), and (4.5) with $\rho < \frac{2\pi}{\sqrt{3}\omega}$.

4. Conclusion

In this paper, we consider a sixth-order variable coefficient singular differential equation with indefinite weights and prove the existence of positive periodic solutions of the equation by the Krasnoselskii-Guo fixed point theorem.

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Appendix

1. Auxiliary Equation

The following are four linear differential equations and the properties of their corresponding Green’s functions.

Equation I

$$\begin{cases} u''(t) + \alpha(t)u'(t) + \beta(t)u(t) = \varphi(t), \\ u(0) = u(\omega), u'(\omega) = u'(0), \end{cases} \tag{4.1}$$

where $\alpha(t), \beta(t) \in C(\mathbb{R}, \mathbb{R}^+)$ are ω - periodic functions. Its solution can be expressed by

$$u(t) = \int_0^\omega G^1(t, s)\varphi(s)ds,$$

where $G^1(t, s)$ is the Green function of equation (4.1).

Lemma 4.1. ([20]) Suppose $f(t) \in X$ and

$$\frac{H[\exp(\int_0^\omega \alpha(x)dx)-1]}{Q\omega} \geq 1, \tag{4.2}$$

where

$$H = \max_{t \in [0, \omega]} \left| \int_0^\omega \frac{\exp(\int_t^s \alpha(x)dx)}{\exp(\int_0^\omega \alpha(x)dx)-1} \beta(s)ds \right|, \quad Q = [1 + \exp(\int_0^\omega \alpha(x)dx)]^2 P^2,$$

Then the equation

$$u''(t) + \alpha(t)u'(t) + \beta(t)u(t) = \varphi(t)$$

has an ω -periodic solution, and its periodic solution can be expressed by

$$u(t) = \int_0^\omega G_1(t, s)\varphi(s)ds,$$

where

$$G^1(t, s) = \frac{\int_t^s \exp[\int_t^x \beta(v)dv + \int_x^s \alpha(v)dv] dx + \int_s^{t+\omega} \exp[\int_t^x \beta(v)dv + \int_x^{s+\omega} \alpha(v)dv] dx}{[\exp(\int_0^\omega \alpha(x)dx) - 1][\exp(\int_0^\omega \beta(x)dx) - 1]}.$$

Lemma 4.2. ([20]) Let

$$A = \int_0^\omega \alpha(x)dx, B = \omega^2 \exp\left(\frac{1}{\omega} \int_0^\omega \ln \beta(x) dx\right).$$

If
$$A^2 \geq 4B, \tag{4.3}$$

then

$$\begin{aligned} \min\{\int_0^\omega \mu(x)dx, \int_0^\omega v(x)dx\} &\geq \frac{1}{2}(A - \sqrt{A^2 - 4B}) := m_1, \\ \max\{\int_0^\omega \mu(x)dx, \int_0^\omega v(x)dx\} &\leq \frac{1}{2}(A + \sqrt{A^2 - 4B}) := n_1. \end{aligned}$$

Lemma 4.3. ([20]) If equations (4.2) and (4.3) hold, then

$$0 < \frac{\omega}{(e^{n_1} - 1)^2} \leq G^1(t, s) \leq \frac{\omega \exp(\int_0^\omega \alpha(x)dx)}{(e^{m_1} - 1)^2}.$$

Equation II

$$\begin{cases} u''(t) + m(t)u(t) = \varphi(t), \\ u(0) = u(\omega), u'(0) = u'(\omega), \end{cases} \tag{4.4}$$

where $m(t) \in C(\mathbb{R}, \mathbb{R}^+)$ is an ω -periodic function. Then, the solution of equation (4.4) can be expressed as

$$u(t) = \int_0^\omega G_2(t, s)\varphi(s)ds,$$

where the Green function $G_2(t, s)$ is defined in [17, Theorem 2.4].

Define

$$k(\beta) = \begin{cases} \frac{2\pi}{\beta \omega^{1+2/\beta}} \left(\frac{2}{2+\beta}\right)^{1-2/\beta} \left(\frac{\Gamma(\frac{1}{\beta})}{\Gamma(\frac{1}{2} + \frac{1}{\beta})}\right)^2, & 1 \leq \beta < +\infty, \\ \frac{4}{\beta}, & \beta = \infty, \end{cases}$$

where $\Gamma(t) = \int_0^{+\infty} u^{t-1} e^{-u} du$.

Lemma 4.4. ([18, Corollary 2.3] and [2, Theorem 3.2]) Assume $m(t) > 0$ and $m(t) \in L^\alpha(0, \omega)$ for some $1 \leq \alpha \leq \infty$. If

$$\|m\| := \left(\int_0^\omega |m(t)|^\alpha dt\right)^{\frac{1}{\alpha}} < k(2\alpha^*), \tag{4.5}$$

where

$$\alpha^* = \begin{cases} \frac{\alpha}{\alpha-1}, & 1 \leq \alpha < +\infty, \\ 1, & \alpha = +\infty, \end{cases}$$

then $G_2(t, s) > 0$ for all $(t, s) \in [0, \omega] \times [0, \omega]$.

Equation III

$$\begin{cases} u''(t) - \rho u(t) + \rho^2 u(t) = \varphi(t), \\ u(0) = u(\omega), u'(0) = u'(\omega), \end{cases} \tag{4.6}$$

where ρ is a constant. Obviously, the solution of equation (4.6) can be expressed as

$$u(t) = \int_0^\omega G_3(t, s)\varphi(s)ds,$$

where the Green’s function

$$G_3(t, s) = \begin{cases} \frac{2e^{\frac{\rho}{2}(t-s)} \left[\sin \frac{\sqrt{3}}{2} \rho(\omega - t + s) + e^{-\frac{\rho\omega}{2}} \sin \frac{\sqrt{3}}{2} \rho(t-s) \right]}{\sqrt{3}\rho(e^{\frac{\rho\omega}{2}} + e^{-\frac{\rho\omega}{2}} - 2 \cos \frac{\sqrt{3}}{2} \rho\omega)}, & 0 \leq s \leq t \leq \omega, \\ \frac{2e^{\frac{\rho}{2}(\omega+t-s)} \left[\sin \frac{\sqrt{3}}{2} \rho(s-t) + e^{-\frac{\rho\omega}{2}} \sin \frac{\sqrt{3}}{2} \rho(\omega-s+t) \right]}{\sqrt{3}\rho(e^{\frac{\rho\omega}{2}} + e^{-\frac{\rho\omega}{2}} - 2 \cos \frac{\sqrt{3}}{2} \rho\omega)}, & 0 \leq t \leq s \leq \omega, \end{cases}$$

Lemma 4.5. ([17]) Assume $\rho < \frac{2\pi}{\sqrt{3}\omega}$ holds, then

$$0 < m_2 := \frac{2 \sin(\frac{\sqrt{3}}{2} \rho\omega)}{\sqrt{3}\rho(e^{\frac{\rho\omega}{2}} + 1)} \leq G_3(t, s) \leq \frac{2}{\sqrt{3}\sin(\frac{\sqrt{3}}{2} \rho\omega)} := n_2, \forall s, t \in [0, \omega].$$

Equation IV

$$\begin{cases} u'(t) + \rho u(t) = \varphi(t), \\ u(0) = u(\omega), \end{cases} \tag{4.7}$$

where ρ is a constant. Similarly, the solution of equation (4.7) can be expressed as

$$u(t) = \int_0^\omega G_4(t, s)\varphi(s)ds,$$

Where

$$G_4(t, s) = \begin{cases} \frac{e^{-\rho(t-s)}}{1 - e^{-\omega\rho}}, & 0 \leq s \leq t \leq \omega, \\ \frac{e^{-\rho(\omega+t-s)}}{1 - e^{-\omega\rho}}, & 0 \leq t \leq s \leq \omega. \end{cases}$$

2. Regularity of the Green function of equation (2.1)

In this section, Cases 2-6 will be presented.

Case 2 There exist two ω -periodic differentiable functions $p(t), q(t): R \rightarrow R^+$ with $\int_0^\omega \alpha(s)ds > 0$ and $\int_0^\omega \beta(s)ds > 0$, and a positive real constant ρ , such that

$$\begin{aligned} a_5(t) &= \alpha(t) - 2\rho, \\ a_4(t) &= 4\alpha'(t) + \beta(t) - 2\rho\beta(t) + \rho^2, \\ a_3(t) &= 5\alpha''(t) + 4\beta'(t) - 6\rho\alpha'(t) - 2\rho\beta(t) + 3\rho^2\alpha(t) - 2\rho^3(t), \\ a_2(t) &= 6m''(t) + 6m'(t)\alpha(t) + m(t)[2\alpha'(t) + 2\beta(t) + \alpha^2(t)] + \beta''(t) + \beta'(t)\alpha(t) + \beta^2(t), \\ a_1(t) &= 4m''(t) + 6m''(t)\alpha(t) + 2m'(t)[2\alpha'(t) + 2\beta(t) + \alpha^2(t)] + m(t)[\alpha''(t) + 2\beta'(t) + 2\beta(t)\alpha(t)], \\ a_0(t) &= m^{(4)}(t) + m'''(t)[2\alpha(t) + \alpha'(t)\alpha(t)] + m''(t)[2\alpha'(t) + 2\beta(t) + \alpha^2(t)] + m'(t)[2\beta'(t) + 2\alpha(t)\beta(t) + \alpha''(t) + \alpha'(t)\alpha(t)] + m(t)[\beta''(t) + \alpha(t)\beta'(t) + \beta^2(t)]. \end{aligned}$$

In this case, equation (2.1) can be transformed into

$$\begin{cases} v''(t) - \rho v'(t) + \rho v(t) = \varphi(t), \\ v(0) = v(\omega), v'(0) = v'(\omega), \end{cases}$$

$$\begin{cases} y''(t) - \rho y'(t) + \rho^2 y(t) = v(t), \\ y(0) = y(\omega), y'(0) = y'(\omega), \end{cases}$$

and

$$\begin{cases} u''(t) + a(t)u'(t) + \beta(t)u(t) = y(t), \\ u(0) = u(\omega), u'(0) = u'(\omega). \end{cases}$$

And the solutions of these second-order equations can be respectively expressed as

$$v(t) = \int_0^\omega G_3(t, s)\varphi(s)ds, y(t) = \int_0^\omega G_3(t, s)v(s)ds, u(t) = \int_0^\omega G_1(t, s)y(s)ds.$$

Similar to Case 1, the solutions of equation (2.1) can be written as

$$u(t) = \int_0^\omega G^2(t, s)\varphi(s)ds,$$

where

$$G^2(t, s) = \int_0^\omega \int_0^\omega G_3(t, \gamma_1) G_3(\gamma_1, \gamma_2) G_1(\gamma_2, s) d\gamma_2 d\gamma_1.$$

Lemma 4.6.[14] Suppose that equations (4.2), (4.3), and $\rho < \frac{2\pi}{\sqrt{3}\omega}$ hold, then $G^2(t, s) > 0$ for any $(t, s) \in [0, \omega] \times [0, \omega]$.

Case 3 There exists a positive constant ρ and an ω -periodic differentiable function $m(t) > 0$ such that

$$\begin{aligned} a_5(t) &= -2\rho, \\ a_4(t) &= m(t) + 3\rho^2, \\ a_3(t) &= 4m'(t) - 2\rho m(t) - 2\rho^3, \\ a_2(t) &= 6m''(t) - 6\rho m'(t) + 3\rho^2 m(t) + \rho^4, \\ a_1(t) &= 4m'''(t) - 6\rho m''(t) + 6\rho^2 m'(t) - 2\rho^3 m(t), \\ a_0(t) &= m^{(4)}(t) - 2\rho m'''(t) + 3\rho^2 m''(t) - 2\rho^3 m'(t) + \rho^4 m(t). \end{aligned}$$

In this case, equation (2.1) can be transformed into

$$\begin{cases} v''(t) - \rho v'(t) + \rho^2 v(t) = \varphi(t), \\ v(0) = v(\omega), v'(0) = v'(\omega), \end{cases}$$

$$\begin{cases} y''(t) - \rho y'(t) + \rho^2 y(t) = v(t), \\ y(0) = y(\omega), y'(0) = y'(\omega), \end{cases}$$

and

$$\begin{cases} u''(t) + m(t)u(t) = y(t), \\ u(0) = u(\omega), u'(\omega) = u'(\omega). \end{cases}$$

And the solutions of these second-order equations can be respectively expressed as

$$v(t) = \int_0^\omega G_3(t, s)\varphi(s)ds, y(t) = \int_0^\omega G_3(t, s)v(s)ds, u(t) = \int_0^\omega G_2(t, s)y(s)ds.$$

Similar to Case 1, the solutions of equation (2.1) can be written as

$$u(t) = \int_0^\omega G^3(t, s)\varphi(s)ds,$$

where

$$G^3(t, s) = \int_0^\omega \int_0^\omega G_3(t, \gamma_1) G_3(\gamma_1, \gamma_2) G_2(\gamma_2, s) d\gamma_2 d\gamma_1.$$

Lemma 4.7. ([14]) Suppose that equation (4.5) and $\rho < \frac{2\pi}{\sqrt{3}\omega}$ hold, then $G^3(t, s) > 0$ for any $(t, s) \in [0, \omega] \times [0, \omega]$.

Case 4 There exist two ω -periodic differentiable functions $\alpha(t), \beta(t): R \rightarrow R^+$ with $\int_0^\omega \alpha(s)ds > 0$ and $\int_0^\omega \beta(s)ds > 0$, and a positive real constant ρ , such that

$$\begin{aligned} a_5(t) &= 2\alpha(t) + 2\rho, \\ a_4(t) &= 4\alpha'(t) + 2\beta(t) + 4\rho\alpha(t) + \alpha^2(t) + \rho^2(t), \\ a_3(t) &= 6\alpha''(t) + 6\rho\alpha'(t) + \alpha(t)[2\rho^2 + 3\alpha'(t) + 2\rho\alpha(t) + 2\beta(t)] + 4\rho\alpha(t) + 4\beta'(t), \\ a_2(t) &= 4\alpha'''(t) + 6\beta''(t) + 2\rho^2\alpha'(t) + 6\alpha\beta'(t) + 6\alpha(t)\alpha''(t) + \rho^2\beta(t) + \beta(t)[2\alpha'(t) + \beta(t) \\ &\quad + 2\rho\alpha(t) + \rho^2] + \alpha(t)[3\alpha''(t) + 3\beta'(t) + 4\rho\alpha'(t) + 2\rho\beta(t) + \rho^2\alpha(t)], \\ a_1(t) &= \rho^2\alpha''(t) + 2\rho^2\beta'(t) + \alpha(t)[3\beta''(t) + 4\rho\beta'(t) + \alpha'''(t) + 2\rho\alpha''(t) + \rho^2\alpha'(t) + \rho^2\beta(t)] \\ &\quad + \beta(t)[\alpha''(t) + 2\rho^2\alpha'(t) + 2\rho\alpha(t) + 2\beta'(t) + \rho^2\beta(t)] + 4\beta'''(t) \\ &\quad + 6\alpha\beta''(t) + 2\rho\alpha''(t) + \alpha^{(4)}(t), \\ a_0(t) &= \beta^{(4)}(t) + 2\rho\beta'''(t) + \alpha^2\beta''(t) + \alpha(t)[\beta''(t) + 2\rho\beta'(t) + \rho^2\beta(t)] + \beta(t)[\beta''(t) \\ &\quad + 2\rho\beta'(t) + \rho^2\beta(t)]. \end{aligned}$$

In this case, equation (2.1) can be transformed into

$$\begin{cases} x''(t) + \alpha(t)x'(t) + p(t)x(t) = \varphi(t), \\ x(0) = x(\omega), x'(\omega) = x'(\omega), \end{cases} \tag{4.8}$$

$$\begin{cases} v'(t) + \rho v(t) = u(t), \\ v(0) = v(\omega), \end{cases} \tag{4.9}$$

$$\begin{cases} y'(t) + \rho y(t) = v(t), \\ y(0) = y(\omega), \end{cases} \tag{4.10}$$

and

$$\begin{cases} u''(t) + \alpha(t)u'(t) + \beta(t)u(t) = y(t), \\ u(0) = u(\omega), u'(\omega) = u'(\omega). \end{cases} \tag{4.11}$$

Since these solutions of equations (4.8) – (4.11) can be respectively expressed as

$$\begin{aligned} x(t) &= \int_0^\omega G_1(t,s)\varphi(s)ds, v(t) = \int_0^\omega G_4(t,s)x(s)ds, \\ y(t) &= \int_0^\omega G_4(t,s)v(s)ds, u(t) = \int_0^\omega G_1(t,s)y(s)ds \end{aligned}$$

Similar to Case 1, the solutions of equation (2.1) can be written as

$$u(t) = \int_0^\omega G^4(t,s)\varphi(s)ds,$$

where

$$G^4(t,s) = \int_0^\omega \int_0^\omega \int_0^\omega G_1(t,\gamma_3) G_4(\gamma_3,\gamma_2)G_4(\gamma_2,\gamma_1)G_1(\gamma_1,s) d\gamma_3d\gamma_2d\gamma_1.$$

Lemma 4.8.([14]) Suppose that equations (4.2) and (4.3) hold, then $G^4(t,s) > 0$ for any $(t,s) \in [0,\omega] \times [0,\omega]$.

Case 5 There exists an ω -periodic differentiable function $m(t) > 0$ and a positive real constant ρ such that

$$\begin{aligned} a_5(t) &= 2\rho, \\ a_4(t) &= 2m(t) + \rho^2, \\ a_3(t) &= 4m'(t) + 4\rho m(t), \\ a_2(t) &= 6m''(t) + 6\rho m'(t) + 2\rho^2m(t) + m^2(t), \\ a_1(t) &= 4m'''(t) + 6\rho m''(t) + 2\rho^2m'(t) + 2m(t)m'(t) + 2\rho m(t), \\ a_0(t) &= m^{(4)}(t) + 2\rho m'''(t) + \rho^2m''(t) + 2\rho m(t)m'(t) + \rho^2m^2(t) + m(t)m''(t). \end{aligned}$$

In this case, equation (2.1) can be transformed into

$$\begin{cases} x''(t) + m(t)x(t) = \varphi(t), \\ x(0) = x(\omega), x'(0) = x'(\omega), \end{cases} \tag{4.12}$$

$$\begin{cases} v'(t) + \rho v(t) = u(t), \\ v(0) = v(\omega), \end{cases} \tag{4.13}$$

$$\begin{cases} y'(t) + \rho y(t) = v(t), \\ y(0) = y(\omega), \end{cases} \tag{4.14}$$

and

$$\begin{cases} u''(t) + m(t)u(t) = y(t), \\ u(0) = u(\omega), u'(0) = u'(\omega). \end{cases} \tag{4.15}$$

Since these solutions of equations (4.12) – (4.15) can be respectively expressed as

$$\begin{aligned} x(t) &= \int_0^\omega G_2(t,s)\varphi(s)ds, v(t) = \int_0^\omega G_4(t,s)u(s)ds, \\ y(t) &= \int_0^\omega G_4(t,s)v(s)ds, u(t) = \int_0^\omega G_2(t,s)y(s)ds. \end{aligned}$$

Similar to Case 1, the solutions of equation (2.1) can be written as

$$u(t) = \int_0^\omega G^5(t,s)\varphi(s)ds,$$

where

$$G^5(t, s) = \int_0^\omega \int_0^\omega \int_0^\omega G_2(t, \gamma_3) G_4(\gamma_3, \gamma_2) G_4(\gamma_2, \gamma_1) G_2(\gamma_1, s) d\gamma_3 d\gamma_2 d\gamma_1.$$

Lemma 4.9.([14]) Suppose that equation (4.5) holds, then $G^5(t, s) > 0$ for any $(t, s) \in [0, \omega] \times [0, \omega]$.

Case 6 There exists a positive real constant ρ such that

$$a_5(t) = 0, a_4(t) = 0, a_3(t) = 2\rho^3, a_2(t) = 0, a_1(t) = 0, a_0(t) = \rho^6.$$

In this case, equation (2.1) can be expressed as

$$\begin{cases} x''(t) - \rho x'(t) + \rho^2 x(t) = \varphi(t), \\ x(0) = x(\omega), x'(0) = x'(\omega), \end{cases} \tag{4.16}$$

$$\begin{cases} v'(t) + \rho v(t) = u(t), \\ v(0) = v(\omega), \end{cases} \tag{4.17}$$

$$\begin{cases} y'(t) + \rho y(t) = v(t), \\ y(0) = y(\omega), \end{cases} \tag{4.18}$$

and

$$\begin{cases} u''(t) - \rho u'(t) + \rho u^2(t) = y(t), \\ u(0) = u(\omega), u'(\omega) = u'(\omega). \end{cases} \tag{4.19}$$

Since these solutions of equations (4.16) – (4.19) can be respectively expressed as

$$\begin{aligned} x(t) &= \int_0^\omega G_3(t, s)\varphi(s)ds, v(t) = \int_0^\omega G_4(t, s)u(s)ds, \\ y(t) &= \int_0^\omega G_4(t, s)v(s)ds, u(t) = \int_0^\omega G_3(t, s)y(s)ds. \end{aligned}$$

Similar to Case 1, the solutions of equation (2.1) can be represented in the form.

$$u(t) = \int_0^\omega G^6(t, s)\varphi(s)ds,$$

where

$$G^5(t, s) = \int_0^\omega \int_0^\omega \int_0^\omega G_3(t, \gamma_3) G_4(\gamma_3, \gamma_2) G_4(\gamma_2, \gamma_1) G_3(\gamma_1, s) d\gamma_3 d\gamma_2 d\gamma_1.$$

Lemma 4.10.([14]) Suppose that $\rho < \frac{2\pi}{\sqrt{3}\omega}$ holds, then $G^6(t, s) > 0$ for any $(t, s) \in [0, \omega] \times [0, \omega]$.